LIGO-Virgo-KAGRA detector and current status of gravitational wave detector

technologies

June Gyu Park



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2. Quantum noise of gravitational wave detector

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Gravitational wave detector and status of LVK





Miller, M.C., Yunes, N. The new frontier of gravitational waves. *Nature* 568, 469–476 (2019)



Gravitational wave



MPA Lectures on Gravitational Waves in Cosmology Azadeh Maleknejad Max-Planck-Institute for Astrophysics





https://www.ligo.caltech.edu/









https://www.ligo.caltech.edu/











Sensitivity of michelson interferometer



Danilishin, Stefan L. et al. Living Rev.Rel. 15 (2012) 5 arXiv:1203.1706



Interferometer of GW detector



































Strain sensitivity





Interferometer of GW detector



IFI : Input faraday isolatorPRM : Power recycling mirrorITM : Input test massETM : End test massMMT : Mode matching telescopeSRM : Signal recycling mirror



PSL room



JGW-G1808402-v6





PSL room



JGW-G1808402-v6



Coherence





Coherence





Linewidth of laser



<u>Transmitter and Receiver Design for Amplified Lightwave Systems</u> Daniel A. Fishman, B. Scott Jackson, in <u>Optical Fiber Telecommunications (Third Edition)</u>, <u>Volume B</u>, 1997



PSL room

Reference Cavity



Status: Installed and automated in air Details:

- A very stable cavity for the frequency reference using the ultra-low expansion (ULE) glass spacer
- Requrement < 100mHz/sec
- Vacuum chamber had a leak; preparing for the repair
- Spacer length 100.71mm (catalog)
- Finesse 30000 (PhD thesis, Nakano)
- Controls: laser thermal, laser PZT, BB EOM
- UGF ~ 500 kHz
- . To be characterized again

JGW-G1808402-v6



Interferometer of GW detector



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Spatial mode of gaussian beam





Gaussian beam



https://www.edmundoptics.co.kr/knowledge-

center/application-notes/lasers/gaussian-beam-propagation/



Gaussian beam





Gaussian beam





Gaussian beam mode matching





Interferometer of GW detector



IFI : Input faraday isolatorPRM : Power recycling mirrorITM : Input test massETM : End test massMMT : Mode matching telescopeSRM : Signal recycling mirror



Gaussian beam mode matching




Interferometer of GW detector



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KAGRA chamber



https://authors.library.caltech.edu/94680/2/1901.03053.pdf



Mirror of IMC





Wavefront sensor MCE_TRANS







Figure 2. Overview of the arm length stabilization system of KAGRA. The frequencies of the two auxiliary lasers are phase-locked to that of PSL. The frequency of the each green laser is controlled by the combination of a double-path acousto-optic

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Vibration isolation system of KAGRA





Cryogenic system of KAGRA



Cryogenic system of KAGRA



Class. Quant. Grav. 36 (2019) 165008



Cryogenic system of KAGRA



T. Akutsu et al. PTEP 2020, 05A101



• Test mass chamber



Rey.Hori



Power and Signal recycling mirror



Y. Aso et al. (KAGRA Collaboration), Phys. Rev. D88, 043007 (2013)

MC : Mode Cleaner ITMX : Input Test Mass X ITMY : Input Test Mass Y REFL : Reflection Port PRM : Power Recycling Mirror OMC : Output Mode Cleaner AS_DC : Anti Symmetric DC SRM : Signal Recycling Mirror POP : Pick-off-in-the-PRC ETMX : End Test Mass X ETMY : End Test Mass Y AS RF : Anti Symmetric RF



Power and Signal-RECYCLED INTERFEROMETER



Power and Signal recycling mirror



Y. Aso et al. (KAGRA Collaboration), Phys.Rev. D88, 043007 (2013)





Power and Signal-RECYCLED INTERFEROMETER



Power and Signal recycling mirror



Y. Aso et al. (KAGRA Collaboration), Phys. Rev. D88, 043007 (2013)

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LIGO status



A+ Upgrade: Toward O4

- □ Target: 190 Mpc BNS range
 - Improved Faraday isolators (UF/Montclair)
 - Adaptive wavefront control (LIGO/Adelaide/Syracuse)
 - High-dynamic-range photodetectors (LIGO/Cardiff)
 - Frequency-dependent squeezing
 - Upgraded squeezer & dark port
 - New end station labs & filter cavity tube enclosure buildings
 - New vacuum chambers & filter cavity beamtubes
 - [+ concurrent detector improvements, not A+: high laser power, point-defect-free TM's, scattered light control, ...]





LIGO Laboratory



A+ Upgrade: Toward O5

- A+ O5 Target: 325 Mpc BNS range
- Low-CTN coatings

LIGO

LIGO-G2200382-v1

- A+ and AdV+ will pursueTi:GeO₂ material for O5; LIGO has procured required deposition targets
 - LIGO/Virgo joint working group appointed to guide full-aperture development & QA
 - Bubble formation risk during annealing is not yet resolved; R&D is underway
- A+ substrates all in hand, now polishing (UK/Glasgow + US)
- Coating production readiness review planned for end of calendar '22
- Large-aperture beamsplitter and relay triple suspension designs complete, fabrication reviews imminent (UK/RAL)
- Balanced Homodyne Readout (BHR) in final design, FDR pending (UK/Glasgow & Cardiff)
 - O5 start paced by post-O4 A+ installation
 - Estimated ~11 months after post-run cal, not including commissioning



BHR raytrace in HAM6, G2101467 (Glasgow)

BBSS Final Design, T2000503 (RAL)



LIGO Laboratory

KAGRA status

Fundamental noises in O4

4



- Observation range limit: 6Mpc
- Target: 1Mpc



- Laser power at BS: 58W
- PRFPMI
- Observation range limit: 35Mpc
- Target: 10Mpc

Noise budget at O3GK



- Noise sources limiting the sensitivity at O3GK were well identified.
 - The noise budget paper has just been submitted to the archive.
- Low frequency (< 100Hz): Suspension control noise.
- Mid-frequency (100Hz-400Hz): Scattering light noise excited by acoustic noise.
- High frequency (400Hz<): Shot noise, optical loss between BS and AS detection port (70% reflection SRM and so on).



Advanced Virgo+ design sensitivity

- Phase I: reduce quantum noise, hit against thermal noise. BNS range: 100 Mpc's
- Phase II: lower the thermal noise wall. BNS range: 200 Mpc's or more





Advanced Virgo+ Phase II



Advanced Virgo

Quantum noise reduction system (QNR)

Goal: use frequency dependent squeezing in AdV+ Phase I

System design compatible with FIS as well



Advanced Virgo+ Phase I



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LIGO status

Updated 16 June 2022	— 01	— 02	— O3	— O4	O 5
LIGO	80 Mpc	100 Мрс	100-140 Мрс	160-190 Mpc	240-325 Mpc
Virgo		30 Mpc	40-50 Мрс	80-115 Mpc	150-260 Mpc
KAGRA			0.7 Mpc	(1-3) ~ 10 Mpc	25-128 Mpc
G2002127-v12 20	 015 2016	1 2017 2018 2	2019 2020 2021	 2022 2023 2024 2025 2	1 1 1 2026 2027 2028





Quantum noise of gravitational wave detector



Quantum noise of coherent light





Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit.Phil. Trans. R. Soc. A 376: 20170289.



Phase and amplitude noise of light



He urs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit.Phil. Trans. R. Soc. A 376: 20170289.



Classical electromagnetic wave







Photo dioe current (I) \propto Intensity \propto (Number of photon)





Photo dioe current (I) \propto Intensity \propto (Number of photon)









Figure 2-3: Balanced homodyne readout.



Quantum noise of coherent light





Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit.Phil. Trans. R. Soc. A 376: 20170289.



Standard quantum limit of GW detector



Standard quantum limit of gravitational wave detector Shot noise + Radiation pressure noise



Quantum noise of interferometer





•Quantum noise of interferometer



Assume anti-symmetric port is dark port

$$E_1(t) = \frac{1}{\sqrt{2}} [E_s(t) + E_{as}(t)]$$
$$E_2(t) = \frac{1}{\sqrt{2}} [E_s(t) - E_{as}(t)]$$

'as' is vacuum field

Squeezed States for Advanced Gravitational Wave Detectors, B.A., University of California Berkeley, Eric Oelker (2009)


Quantum noise of coherent light





Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit.Phil. Trans. R. Soc. A 376: 20170289.







Photon Counting Statistics

















Target sensitivity of KAGRA





Target sensitivity of KAGRA





Shot noise curve of GW detector



Interaction length ~ Number of round trip x Length



Radiation pressure noise



- Stored energy is very high (750 kW)
- Desired sensitivity is very high ($10^{-21} \sim 10^{-24}$)



• Test mass of KAGRA





Radiation pressure noise



$$E_{as,r}(t) = E_{as} + E_0 \frac{\omega_0 [x_2(t) - x_1(t)]}{c} \sin(\omega_0 t)$$

$$t) - x_2(t) = \frac{x_{cl,1}(t) - x_{cl,2}(t) + \delta \hat{x}_1(t) - \delta \hat{x}_2(t) + Lh(t)}{\text{Thermal, seismic}}$$
Radiation pressure GW source (slow)



Design sensitivity of KAGRA





Squeezed vacuum injection in GW detector



•Quantum noise of interferometer



Assume anti-symmetric port is dark port

$$E_1(t) = \frac{1}{\sqrt{2}} [E_s(t) + E_{as}(t)]$$
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Non-linear crystal





Parametric down conversion





Parametric down conversion process





Parametric down conversion





Parametric down conversion process





Squeezed vacuum



Figure 1.6: Simulation of electric field in time for (a) vacuum state and for (b) squeezed vacuum.



Phase and amplitude noise of light



Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit.Phil. Trans. R. Soc. A 376: 20170289.



Squeezed state of light



R. Schnabel / Physics Reports 684 (2017) 1-51



Frequency independent squeezing



Optical and noise studies for Advanced Virgo and filter cavities for quantum noise reduction in gravitational-wave interferometric detectors, Eleonora Capocasa, UNIVERSITÉ PARIS DIDEROT (2017)



Frequency independent squeezing



Optical and noise studies for Advanced Virgo and filter cavities for quantum noise reduction in gravitational-wave interferometric detectors, Eleonora Capocasa, UNIVERSITÉ PARIS DIDEROT (2017)



Frequency independent squeezing



Optical and noise studies for Advanced Virgo and filter cavities for quantum noise reduction in gravitational-wave interferometric detectors, Eleonora Capocasa, UNIVERSITÉ PARIS DIDEROT (2017)



Frequency dependent squeezing(FDS)



Optical and noise studies for Advanced Virgo and filter cavities for quantum noise reduction in gravitational-wave interferometric detectors, Eleonora Capocasa, UNIVERSITÉ PARIS DIDEROT (2017)





4. Frequency dependent squeezing in GW detector



First suggestion of filter cavity in FD squeezing

Conversion of conventional gravitational-wave interferometers into quantum nondemolition interferometers by modifying their input and/or output optics

H. J. Kimble,¹ Yuri Levin,^{2,*} Andrey B. Matsko,³ Kip S. Thorne,² and Sergey P. Vyatchanin⁴ ¹Norman Bridge Laboratory of Physics 12-33, California Institute of Technology, Pasadena, California 91125 ²Theoretical Astrophysics, California Institute of Technology, Pasadena, California 91125 ³Department of Physics, Texas A&M University, College Station, Texas 77843-4242 ⁴Physics Faculty, Moscow State University, Moscow, 119899, Russia



FIG. 1. Schematic diagram of a squeezed-input interferometer.



Squeezed vacuum injection with filter cavity



Side band figure





Detuned cavity

Stefan Hild et al, "Detuned arm cavities", 3rd GEO simulation workshop, Hannover, June 2007

<u>B:</u>

less carrier light in cavity => less GW sidebands are produced.
Since one GW sideband is resonant, it gets enhanced.

Simple picture

=> Smaller GW signal

<u>C:</u>

optical power is restored in the cavity by larger PR-gain.
Same amount of GW sidebands are produced.
Since one GW sideband is

resonant, it gets enhanced. Overall we win GW signal.

=> Larger GW signal



brea Astronomy and

3rd GEO simulation workshop, Hannover, June 2007 pace Science Institute

Gravitational wave signal



Electric field in simple Michelson interferometer



Assume anti-symmetric port is dark port

$$E_1(t) = \frac{1}{\sqrt{2}} [E_s(t) + E_{as}(t)]$$
$$E_2(t) = \frac{1}{\sqrt{2}} [E_s(t) - E_{as}(t)]$$

'as' is vacuum field

Squeezed States for Advanced Gravitational Wave Detectors, B.A., University of California Berkeley, Eric Oelker (2009)



Optical parametric oscillator



All-Optical Electron Acceleration with Ultrafast THz Pulses, Wenqian Ronny Huang, MIT(2017)


Filter cavity



Denis Martynov et al, Phys. Rev. D **99**, 102004





°ω

Detuned filter cavity



P. Kwee, J. Miller, T. Isogai, L. Barsotti, and M. Evans Phys. Rev. D 90, 062006 – Published 5 September 2014





M. Evans, L. Barsotti, P. Kwee, J. Harms, and H. Miao Phys. Rev. D **88**, 022002 – Published 29 July 2013





UNIVERSITÉ PARIS DIDEROT, Eleonora Capocasa (2017)







UNIVERSITÉ PARIS DIDEROT, Eleonora Capocasa (2017)







UNIVERSITÉ PARIS DIDEROT, Eleonora Capocasa (2017)





Capocasa (2017)



$$\alpha_p = \arctan\left(\frac{2\gamma_{\rm fc}\Delta\omega_{\rm fc}}{\gamma_{\rm fc}^2 - \Delta\omega_{\rm fc}^2 + \Omega^2}\right)$$

 $\gamma = loss of filter cavity$ $\omega_{fc} = detuned frequency$

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$$t_{\rm st} = \frac{1}{\gamma_{\rm fc}} = \frac{\sqrt{2}}{\Omega_{
m SQL}} \simeq 3 \,{
m ms}$$

Squeeze angle rotation



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LIGO filter cavity



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• OPO of LIGO squeezer





LIGO OPO





LIGO filter cavity



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ALS with green laser



Fig. 2. (Color online) Schematic of the arm-length stabilisation system. The numbering indicates the flow of the lock acquisition process and corresponds to the enumerated list below.

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LIGO filter cavity



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