# **Results: Timing Properties**





- 1. A transition layer is seen to be forming to accommodate the change in angular momentum at the NS surface.
- 2. When viscosity is raised, this transition layer grows radially.
- 3. A Keplerial disk forms between the transition layer and the sub-Keplerian flow.





# **Results: Spectral Properties**





Bhattacharjee A., Chakrabarti S. K., MNRAS 472, 1361-1371 (2017)

### TCAF around **Neutron Stars**



## **Results: Spectral Properties**

Nata=0033, hv=1.95 keV, 0=147.362"



#### Table 2 List of parameters for the system used in simulation

ID	$\dot{m}_d (M_{EDD})$	mh (MEDD)	$X_s(r_S)$	$T_{CE}$ (keV)	nee (×1018)	$T_{NS}$ (keV)	$\theta_*$ (degrees)	$T_e(\tau_0)$ (keV)
C1	0.1	0.1	46.8	25.0	1.772	0.802	11.52	22.062
C2	0.1	0.2	46.8	25.0	2.658	0.888	17.48	22.062
C3	0.1	0.5	46.8	25.0	5.317	1.056	36.90	22.062
C4	0.2	0.1	30.0	10.0	1.515	0.888	17.48	8.909
C5	0.2	0.2	30.0	10.0	2.020	0.954	23.61	8.909
C6	0.2	0.5	30.0	10.0	3.535	1.098	44.40	8.909
C7	0.5	0.1	21.8	3.0	2.100	1.056	36.90	2.705
CB	0.5	0.2	21.8	3.0	2.450	1.098	44.40	2.705
C9	0.5	0.5	21.8	3.0	3.500	1.200	90.00	2.705

Bhattacharjee A., Chakrabarti S. K., MNRAS 472, 1361-1371 (2017)

Num=0079, hv=121.83 keV, 0=121.657°

N<sub>arat</sub>=0088, hv=17.31 keV, 0=76.9176\*

### **TCAF** around **Neutron Stars**





- 1. Both the disk and halo accretion rate control the spectral shape.
- Increase in halo 2. accretion increases the hardness of the spectra.
- The simulated spectra 3. matches well with multiple neutron stars Xray binaries.



# **Results: Spectral Properties**

#### (b) $\dot{m}_{d}$ =0.1, $\dot{m}_{h}$ =0.5, $T_{e}(\tau_{0}/2)^{NS}$ =4.5 keV, $T_{e}(\tau_{0}/2)^{KD}$ =41.5 keV



(d)  $\dot{m}_{d}$ =0.1,  $\dot{m}_{h}$ =0.9,  $T_{e}(\tau_{0}/2)^{NS}$ =7.4 keV,  $T_{e}(\tau_{0}/2)^{KD}$ =35.1 keV



#### (a) $m_d=0.1$ , $m_h=0.3$ , $T_e(\tau_0/2)^{NS}=7.2 \text{ keV}$ , $T_e(\tau_0/2)^{KD}=64.7 \text{ keV}$



(c)  $\dot{m}_{d}$ =0.1,  $\dot{m}_{h}$ =0.7,  $T_{e}(\tau_{0}/2)^{NS}$ =3.6 keV,  $T_{e}(\tau_{0}/2)^{KD}$ =30.8 keV



Bhattacharjee A., Chakrabarti S. K., MNRAS 472, 1361-1371 (2017)

-0.45 0.4 0.35 Disk Black-body 0.3 0.25 0.2 0.15 0.05

#### TCAF around **Neutron Stars**

- Cooling of CENBOL. 1.
- Shift of peak temperature location. 2.
- 3. Matches with models used for spectral analysis.





### Tcaf Around Compact Objects or... TACO??

### Schematic of the Model



For BH:  $T_{CO} = 0.0$ ,  $f_{IC} = 0.0$ ,  $f_{CB} = 0.0$ 



## **Use of TCAF: Case Studies**



Bhattacharjee et al. 2017, MNRAS, 466, 1372



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Banerjee et al. 2019, arXiv, 1904.11644

## **TCAF Modelling of BH Sources**

H1743-322



Bhattacharjee et al. 2017, MNRAS, 466, 1372



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Banerjee et al. 2019, arXiv, 1904.11644

# **TACO: Simulated Spectra of BH**

Variation with  $M_{BH}$ 



 $M_{BH} = 5.0,$  $M_{BH} = 10.0,$  $M_{BH} = 15.0$ 



# **TACO: Simulated Spectra of NS**



## **Conclusions: P1**

Accretion around **Compact Objects** 

- temporal variation of black holes and neutron stars, especially in the X-ray (Bhattacharjee 2018).
- Simulating accretion around such objects can be done with a combination of radiation loss within a hydrodynamic code (Bhattacharjee and Chakrabarti 2019).
- Chakrabarti 2020, 2021).
- Jets and outflows are produced due to the formation of shocks in the accretion (Bhattacharjee and Chakrabarti 2019).



• The accretion process around compact objects can be explained with two components of accretion: a Keplrian disk and a sub-Keplerian disk (Chkarabarti and Titarchuk 1995). This framework seems to provide a reasonable insight into the spectral and

• Formation of Keplerian disk requires a treatment for viscosity and more realistic treatment of radiative loss (Bhattacharjee and

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# Numerical Simulations of Accretion-Ejection around Compact Objects: What to include (and what not to)? P2. Study of Jets

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# What are Radio Galaxy Jets?



Image Credits: Frank M. Reiger



### **Morphological Dichotomy of Jets: Fanaroff-Riley Classification**





Radio galaxy Cygnus A at 5 GHz , d ~ 220 Mpc, (z=0.056), extension ~120 kpc (credits: NRAO/AUI, A. Bridle)





Image Credits: Frank M. Reiger



### What are the key characteristics of FR-I?



Hardcastle and Croston, 2011, MNRAS, 415, 133

### Diffused emission in radio

### Deceleration of jet





## **Observational Outlook: Evolving Dichotomy**









# Ledlow and Owen 1996

### Mingo et al. 2019



# **Observational Outlook: Present View**

#### Low-Powered FR-II



### **Overlap of FR-I/II**



FR-I

0





(c) FRII, all structures connected



Mingo et al. 2019

(d) FRII, structures not connected

## **Theoretical Outlook: Present View**





### CFI





KHI

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#### Perucho+ 2010

Gourgouliatos, Komissarov 2018

# **Observation vs Theory**





#### Laing and Bridle 2002



Laing and Bridle 2014, many more

# **Simulation Setup**

#### **King Profile**

$$P_{ext} = \frac{k_B T_{ext}}{\mu X} n_{ext} = \frac{k_B T_{ext}}{\mu X} n_0 \Big[ 1 + \Big(\frac{r}{r_c}\Big)^2 \Big]^{-3\beta_{atm}/2};$$

 $\mu = 0.5, X = 1, n_0 = 0.18 cm^{-3}, r_c = 1.2 kpc, \beta_{atm} = 0.73.$ 

Perucho et al. 2014, MNRAS, 441, 1488

#### **Power-Law Profile**

 $P_{ext} = (5.7 \pm 0.9) \times 10^{-11} (r/r_0)^{-1.5 \pm 0.2} dyn \ cm^{-2};$  $r_0 = 1.0 \ kpc$ .

Wykes et al. 2019, MNRAS, 485, 872

Stellar Mass-Loading Rate: 
$$Q = Q_0 \left(\frac{r_b}{r}\right)^{\gamma} \left[1 + \frac{r_b}{r}\right]^{\gamma} \left[1 + \frac{r_b$$



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Perucho et al. 2014, MNRAS, 441, 1488

### **Fiducial Set: Variation of Flow Structure with Jet Power**



Bhattacharjee, Seo, Ryu, Kang, 2023a (in prep)



# **Simulating Radio Maps**

 $D_{\rm L} = 1/(\Gamma[1-\beta\cos\theta_{\rm obs}])$ 

$$\begin{split} j_{\nu'}^{'} &= \frac{4}{9} \Big( \frac{q^2}{mc^2} \Big)^2 u_B^{'} \nu^{'\frac{1}{2}} \nu_0^{'-\frac{3}{2}} \beta_e^{'2} c \mathcal{N}_e^{'} \Big( \sqrt{\frac{\nu'}{\nu_0'}} \Big), \\ j_{\nu'}^{'} &= -\frac{2}{9} \frac{\mathcal{P} + 2}{m\nu'^2} \Big( \frac{q^2}{mc^2} \Big)^2 u_B^{'} \nu_0^{'-1} \beta_e^{'2} c \mathcal{N}_e^{'} \Big( \sqrt{\frac{\nu'}{\nu_0'}} \Big). \end{split}$$

$$\begin{split} j_{\nu'}^{'} &= \frac{4}{9} \Big( \frac{q^2}{mc^2} \Big)^2 u_B^{'} \nu^{'\frac{1}{2}} \nu_0^{'-\frac{3}{2}} \beta_e^{'2} c \mathcal{N}_e^{'} \Big( \sqrt{\frac{\nu'}{\nu_0'}} \Big), \\ \alpha_{\nu'}^{'} &= -\frac{2}{9} \frac{\mathcal{P}+2}{m\nu'^2} \Big( \frac{q^2}{mc^2} \Big)^2 u_B^{'} \nu_0^{'-1} \beta_e^{'2} c \mathcal{N}_e^{'} \Big( \sqrt{\frac{\nu'}{\nu_0'}} \Big). \end{split}$$

$$j_{\nu} = D_{\rm L}^2 j_{\nu'}' (\nu_{\rm obs}/D_{\rm L}),$$
  
 $\alpha_{\nu} = D_{\rm L}^{-1} \alpha_{\nu'}' (\nu_{\rm obs}/D_{\rm L}).$ 

$$\alpha_{\nu} = L$$

 $dI_{\nu}$ ds

$$\nu_{obs}$$

### Estimating $B_{comov}$

$$B_{p}^{2} = \frac{p}{\beta}; \beta = 100.$$

$$B_{turb}^{2} \approx E_{K,turb}.$$

$$B_{Bell}^{2} \approx \frac{3v_{s}}{2c}(0.1\rho_{1}v_{s}^{2}).$$

$$B_{comov} = \max(B_{p}, B_{turb}, B_{Bell})$$

$$= j_{\nu} - \alpha_{\nu} I_{\nu}.$$

= 1.5 GHz

Estimating  $\mathcal{N}_{e}^{'}$ 

$$\mathcal{N}_{e}^{'}(\gamma^{'})=\mathcal{N}_{0}^{'}\gamma^{'-\mathcal{P}}.$$

$$\mathcal{N}_{0}^{'} = \left(\frac{e^{'}(\mathcal{P}-2)}{1-C_{E}^{'2-\mathcal{P}}}\right)^{\mathcal{P}-1} \left(\frac{1-C_{E}^{'1-\mathcal{P}}}{\frac{\rho^{'}}{m_{p}}(\mathcal{P}-1)}\right)^{\mathcal{P}-2}$$

$$p = 3.0, C'_E = 10^3.$$

Bhattacharjee, Seo, Ryu, Kang, 2023a (in prep)





## Fiducial Set: Variation of Radio Image with Jet Power



Bhattacharjee, Seo, Ryu, Kang, 2023a (in prep)



### **Effects of Jet Density**





Bhattacharjee, Seo, Ryu, Kang, 2023a (in prep)



### Effects of Scales of Propagation and Jet Radius





Bhattacharjee, Seo, Ryu, Kang, 2023a (in prep)



### Effects of Ambient Media: Steeper Background vs Background with Mass-Loading





Bhattacharjee, Seo, Ryu, Kang, 2023a (in prep)

![](_page_25_Picture_4.jpeg)

### **Effects of Ambient Media: King vs Power-Law**

![](_page_26_Figure_1.jpeg)

2023b (in prep)

![](_page_26_Picture_4.jpeg)

### **Effects of Ambient Media: Mass-Loading**

![](_page_27_Figure_1.jpeg)

Bhattacharjee, Seo, Ryu, Kang, 2023b (in prep)

![](_page_27_Picture_3.jpeg)

![](_page_27_Picture_4.jpeg)

# Effects of Ambient Media: $Q_j \sim 10^{42} - 10^{45}$ erg/s

![](_page_28_Figure_1.jpeg)

Bhattacharjee, Seo, Ryu, Kang, 2023b (in prep)

![](_page_28_Picture_3.jpeg)

![](_page_28_Picture_4.jpeg)

# Effects of Mass-Loading: $Q_i \sim 10^{42} - 10^{45}$ erg/s

![](_page_29_Figure_1.jpeg)

![](_page_29_Picture_2.jpeg)

## **Conclusions: P2**

- The jets ejected from AGN requires different treatment at different stages of their lifetime.
- Expanded jets, which have reached ICM and have reached mildly relativistic speeds can be treated well within a purely hydrodynamic setup (Massaglia et el. 2016).
- Launching of jets do not require strong magnetic fields and is most likely can have hydrodynamic origin.
- treatment requires further comparative studies.

![](_page_30_Picture_7.jpeg)

Jets around AGN

• Jets that transition from ultra-relativistic to mildly relativistic or non-relativistic speeds requires Relativistic Hydrodynamic simulations for correctly capturing their dynamics [Rossi et al. 2008; Perucho et al. 2014; Bhattacharjee et al. 2023 (in prep)].

• Magnetic fields can be important in the acceleration phase of jets, however, radiative acceleration can also be important and this

![](_page_30_Picture_15.jpeg)