

Stochastic Gravitational Waves and Its Detection

Chan Park (IBS)

2023.06.14 @ APCTP

The 69th Workshop on Gravitational Waves and Numerical Relativity

Motivations

- The discovery of stochastic gravitational waves background (SGWB) from supermassive binaries is just around the corner.
- The second wave of GW astronomy is coming.
- I expect that SGWB will be the basic course of GW science.
- I will give an intensive basic lecture of SGWB for astrophysicists including the properties of stochastic GWs and working principles of its detectors.

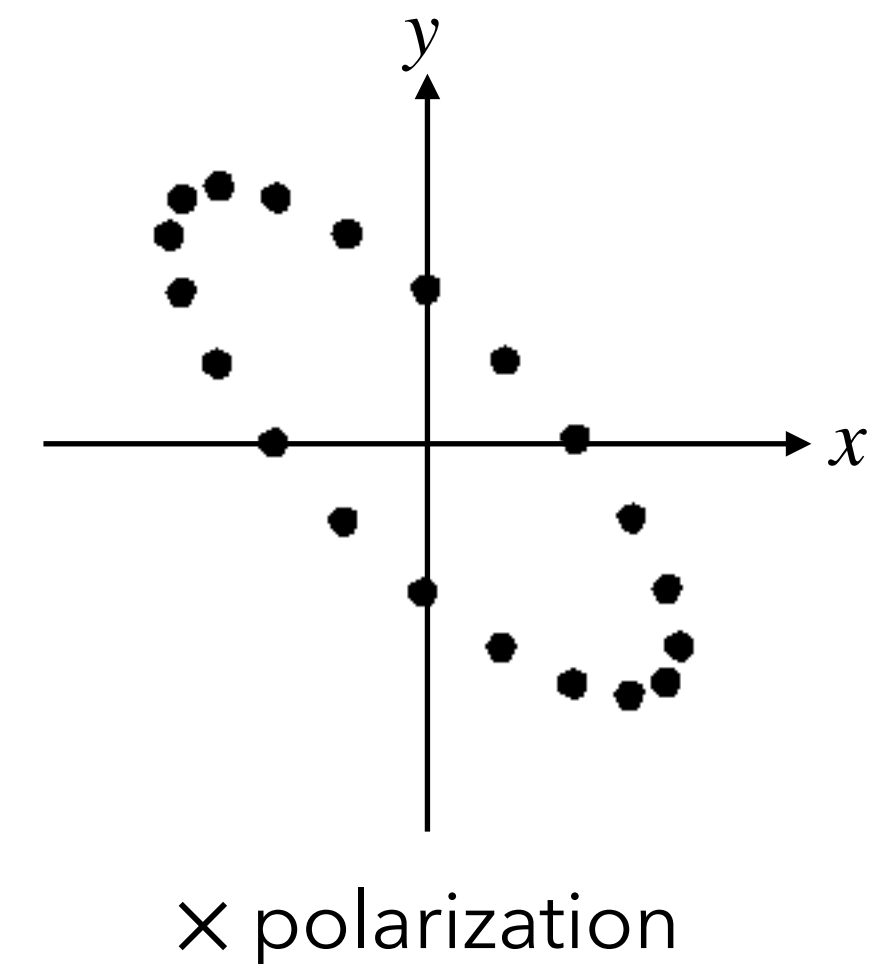
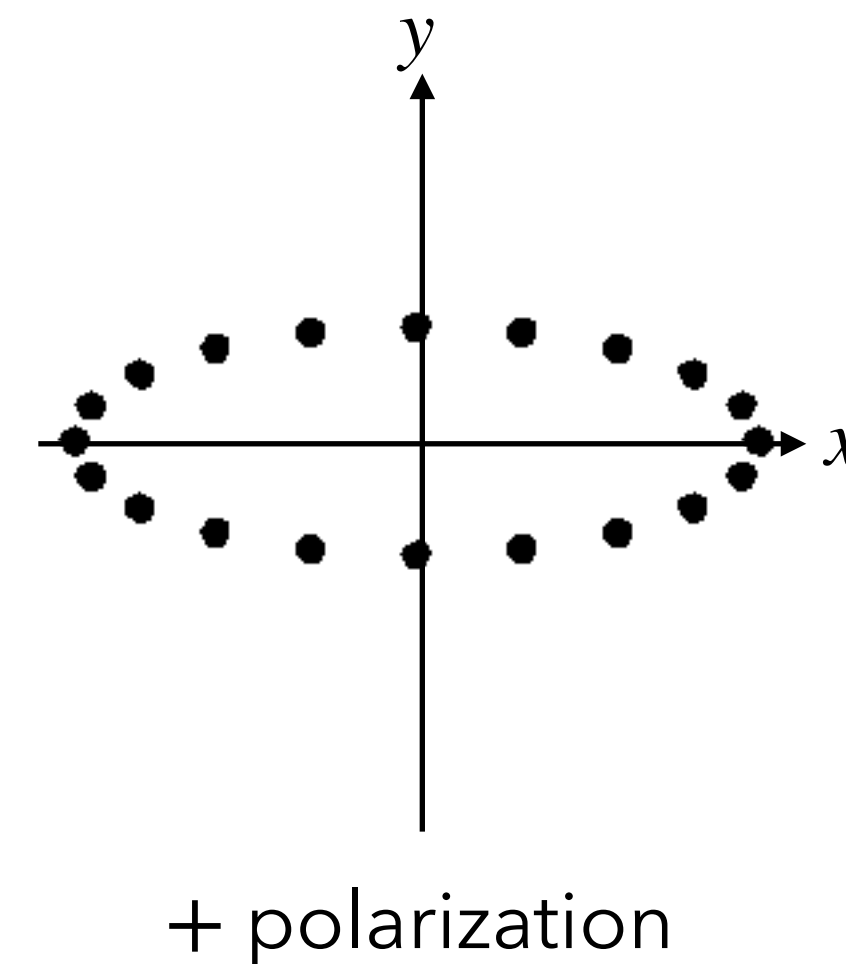
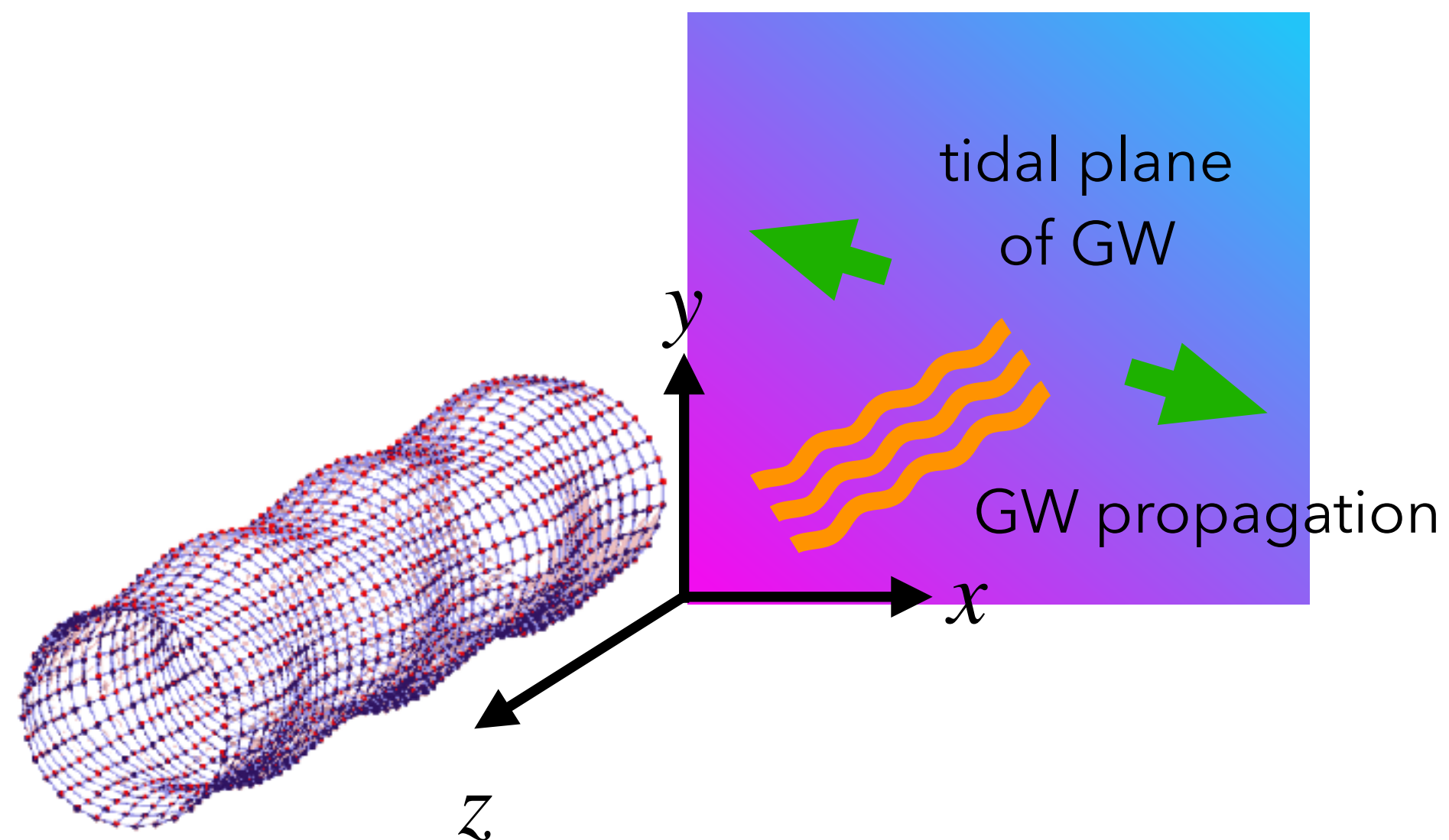
Overview

- Stochastic Gravitational Waves (SGWs)
- Noise Reduction by Correlation Method
- Pulsar Timing Arrays (PTAs)
- Detection of SGWs by Electromagnetic Cavities



Properties of Gravitational Waves (GWs)

- Propagation speed: speed of light
- Transverse-Traceless Gauge
 - Transverse wave: propagation direction $\perp$ tidal direction
 - No expansion of tidal plane: GWs do not change the area, but the shape.
 - Two polarization modes: plus polarization and cross polarization

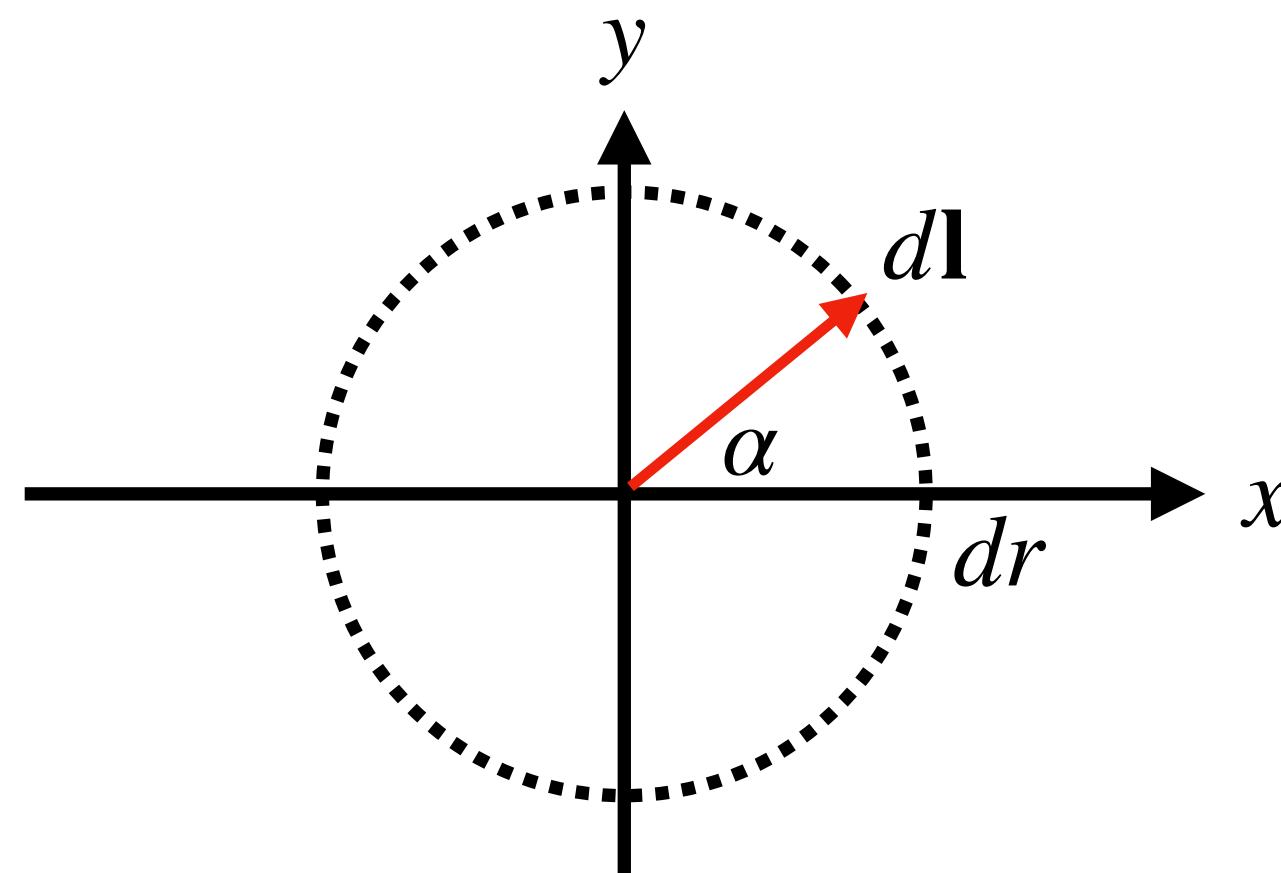


Metric Perturbation of GWs

- Let us consider a metric perturbation $h_{\mu\nu} \ll 1$ preserving area of tidal plane

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & h_+ & h_\times & 0 \\ 0 & h_\times & -h_+ & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

- Let us consider the coordinate displacement $d\mathbf{l} = (0, dr \cos \alpha, dr \sin \alpha, 0)$

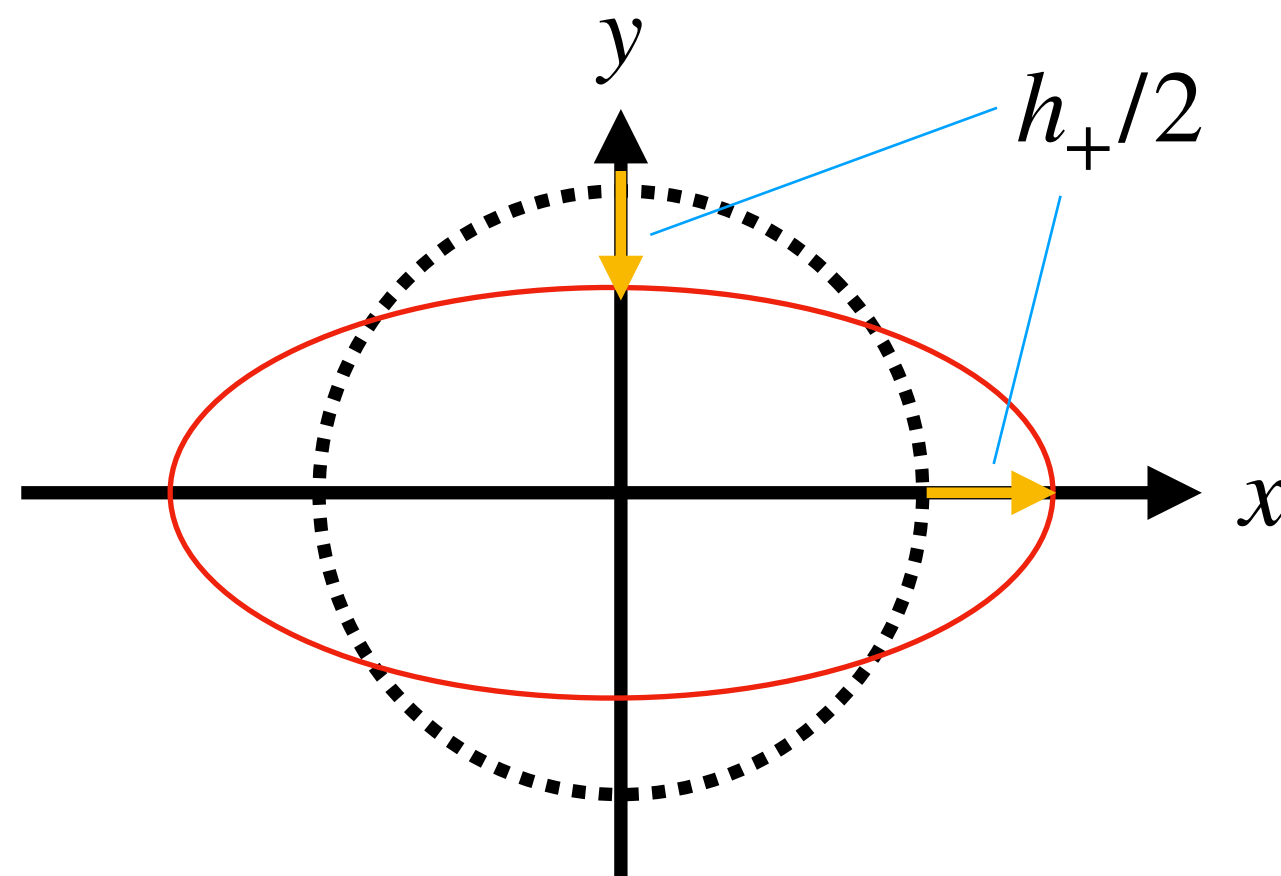


Metric Perturbation of GWs

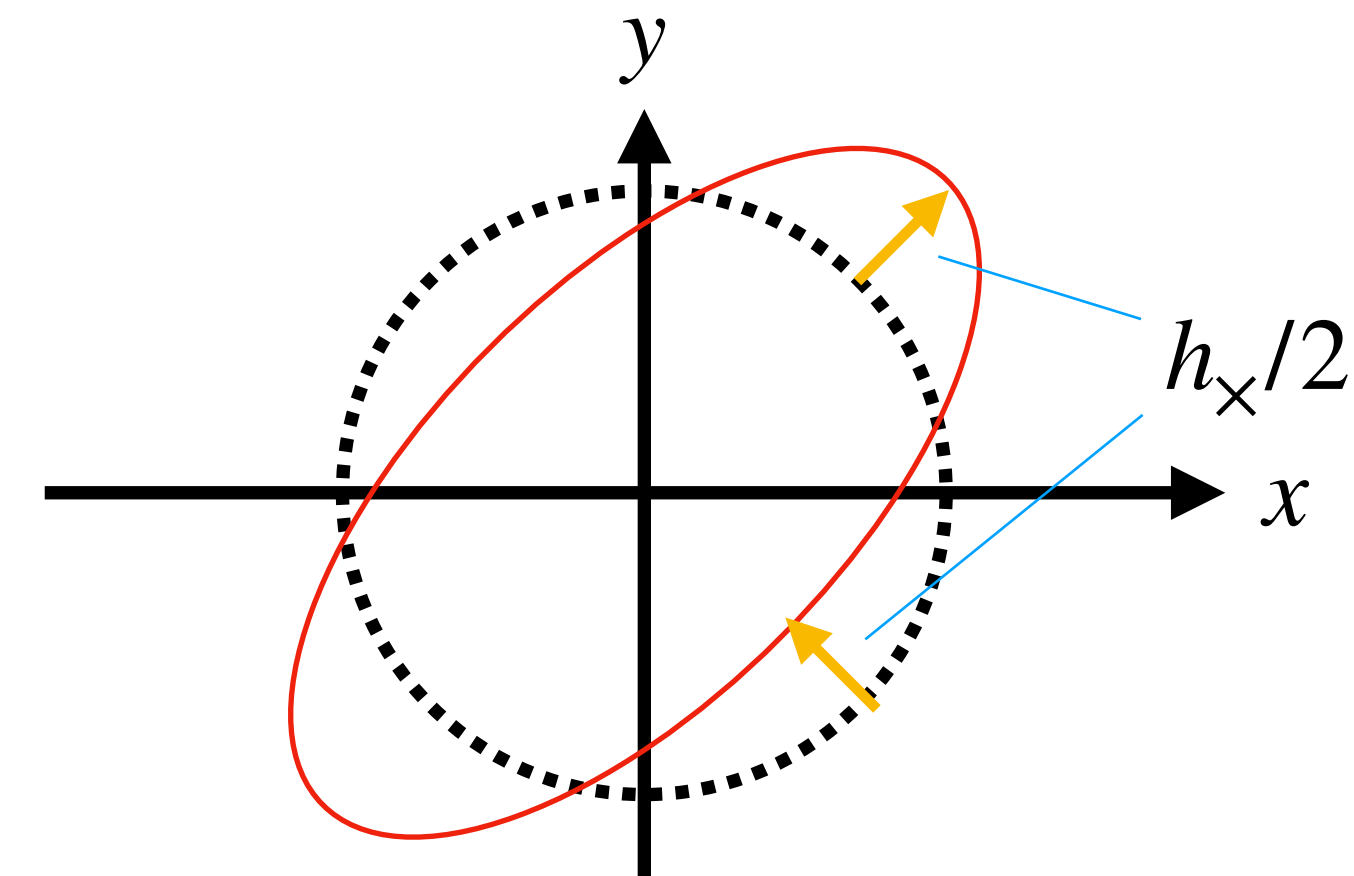
- The physical length of $d\mathbf{l}$ is given by

$$\bullet \quad d\mathbf{l} \cdot d\mathbf{l} = [0 \quad dr \cos \alpha \quad dr \sin \alpha \quad 0] \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 + h_+ & h_\times & 0 \\ 0 & h_\times & 1 - h_+ & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ dr \cos \alpha \\ dr \sin \alpha \\ 0 \end{bmatrix}$$

$$\bullet \quad \sqrt{d\mathbf{l} \cdot d\mathbf{l}} = dr \left\{ 1 + \frac{1}{2} h_+ \cos(2\alpha) + \frac{1}{2} h_\times \sin(2\alpha) \right\}$$



+ polarization

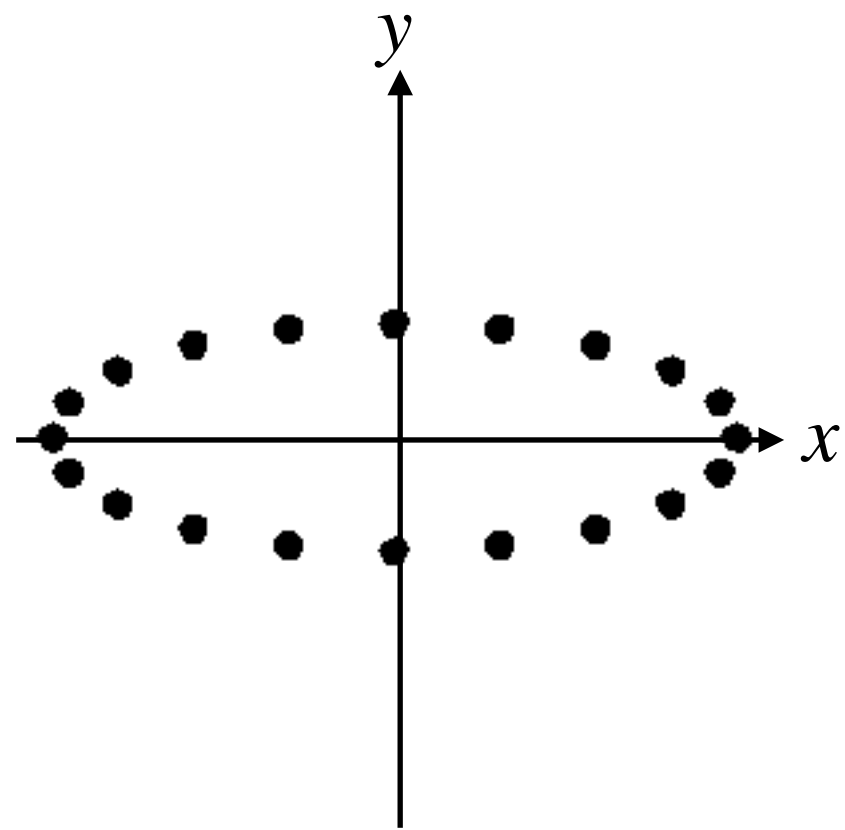


x polarization

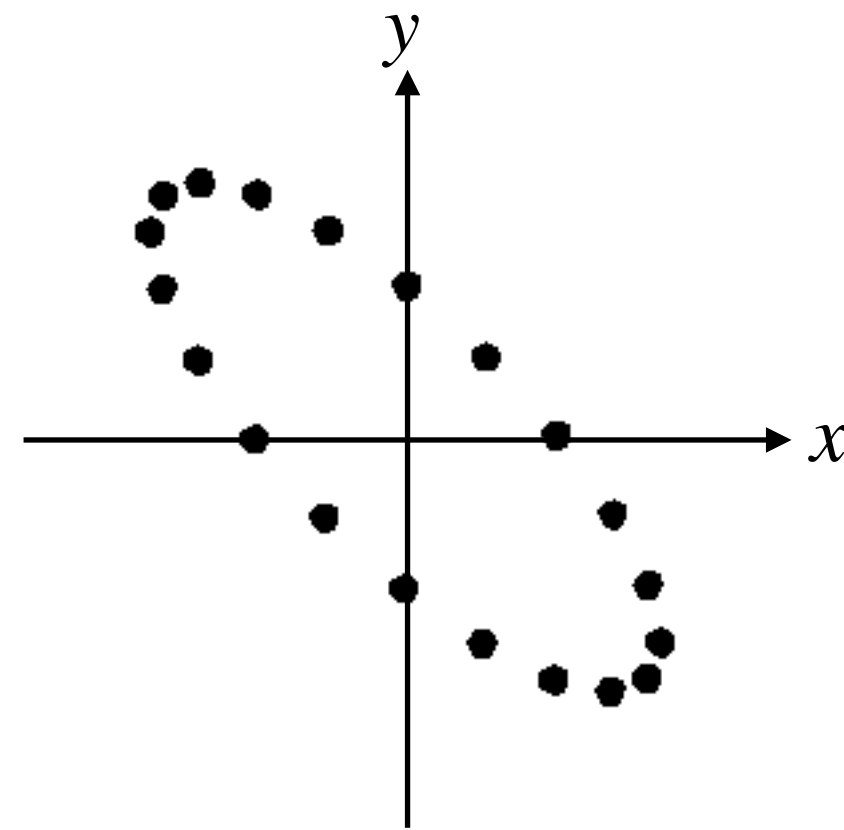
Gravitational Waves

- Monochromatic Plane GWs propagating to $+z$ axis

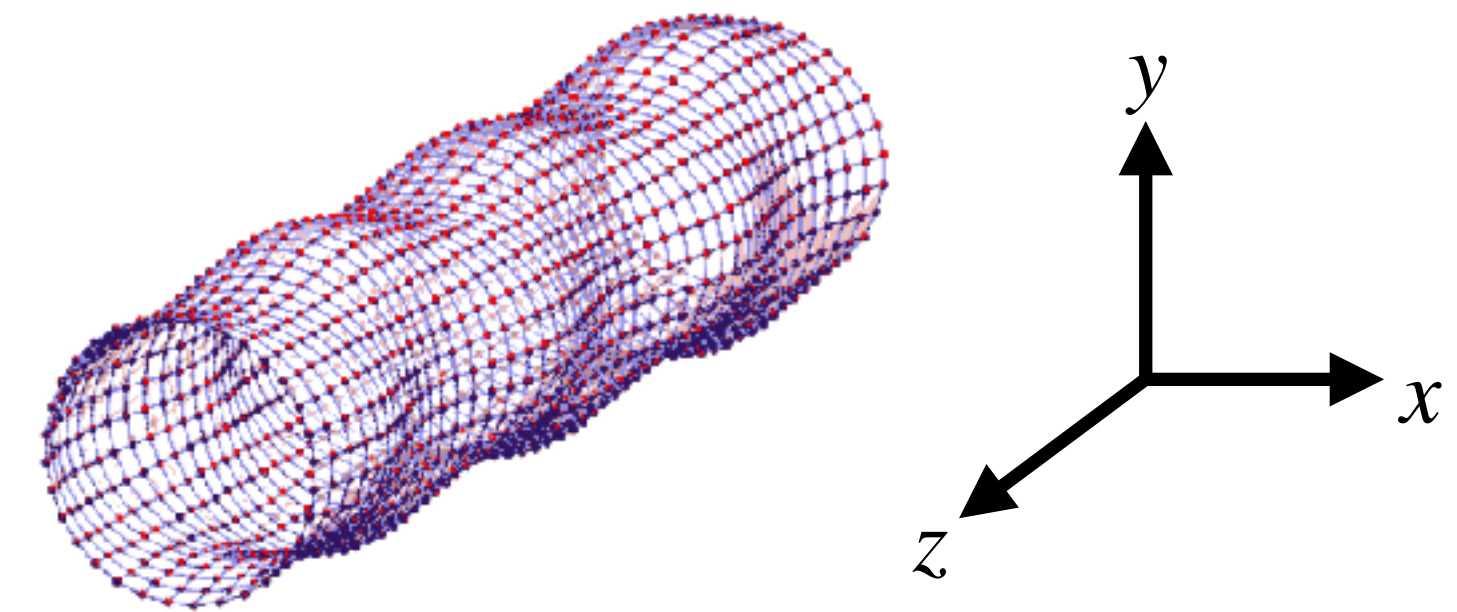
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & h_+ & h_\times & 0 \\ 0 & h_\times & -h_+ & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \cos \left\{ \omega_g (-t + z) + \phi \right\}$$



+ polarization



x polarization

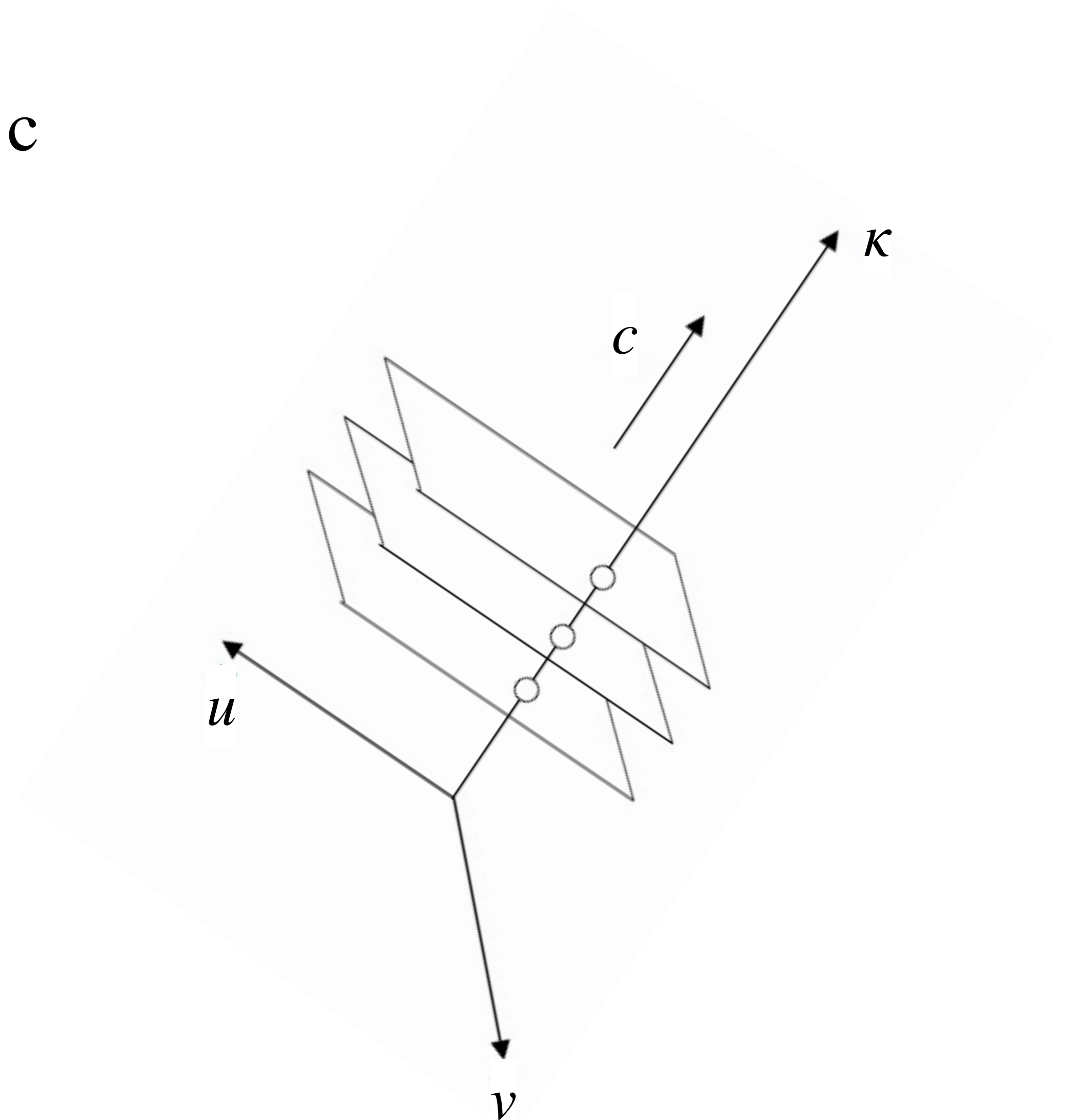


propagation of GWs

Stochastic Gravitational Waves

Monochromatic Plane GWs

- $h_{ab}(t, \vec{x}) = 2\Re \left[\tilde{h}_{ab} e^{iP(t, \vec{x})} \right] = \tilde{h}_{ab} e^{iP(t, \vec{x})} + \text{c.c}$
- Complex amplitude: $\tilde{h}_{ab}$
- Phase: $P(t, \vec{x}) = \omega(-t + \kappa \cdot \vec{x})$
- Angular Frequency: ω
- Unit vector of propagation direction: κ
- Transverse-Traceless Gauge condition
 - $\tilde{h}^a_a = 0$ $\tilde{h}_{ab} n^b = 0$ $\tilde{h}_{ab} \kappa^b = 0$

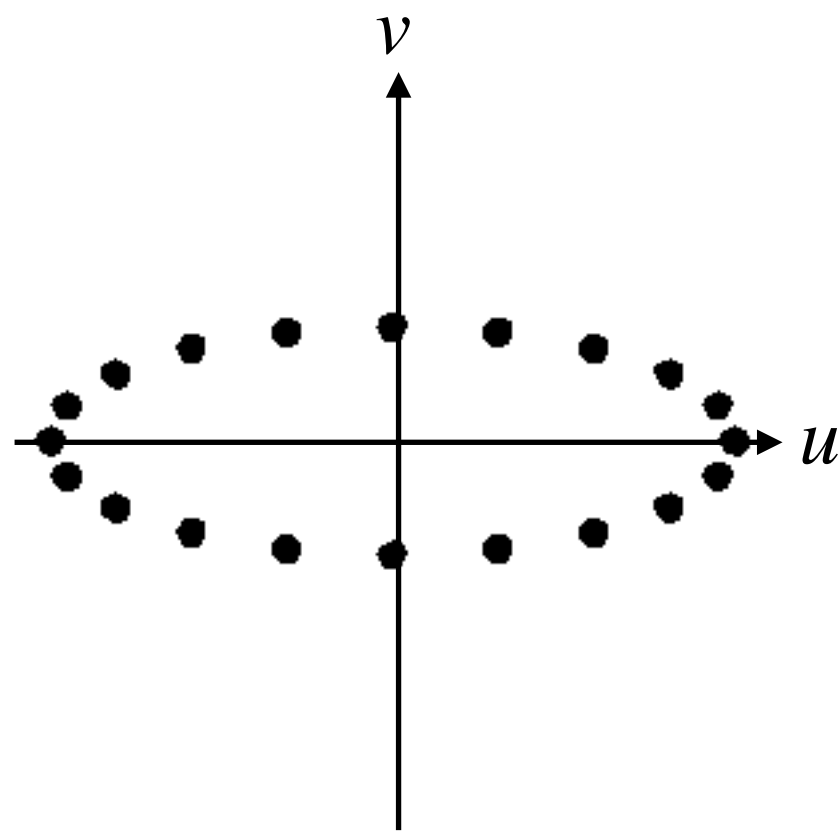


Polarization of GWs

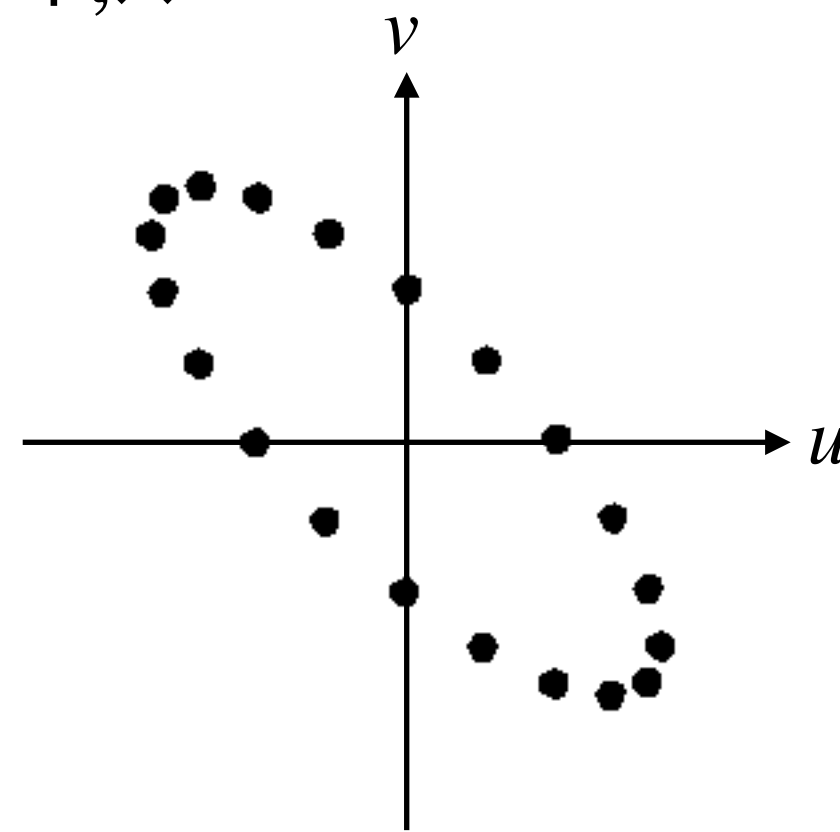
- $e_{ab}^+ = \frac{1}{\sqrt{2}} (u_a u_b - v_a v_b)$ s.t. $e_{ab}^A e_{cd}^B g^{ac} g^{bd} = \delta^{AB}$

- $e_{ab}^\times = \frac{1}{\sqrt{2}} (u_a v_b + v_a u_b)$

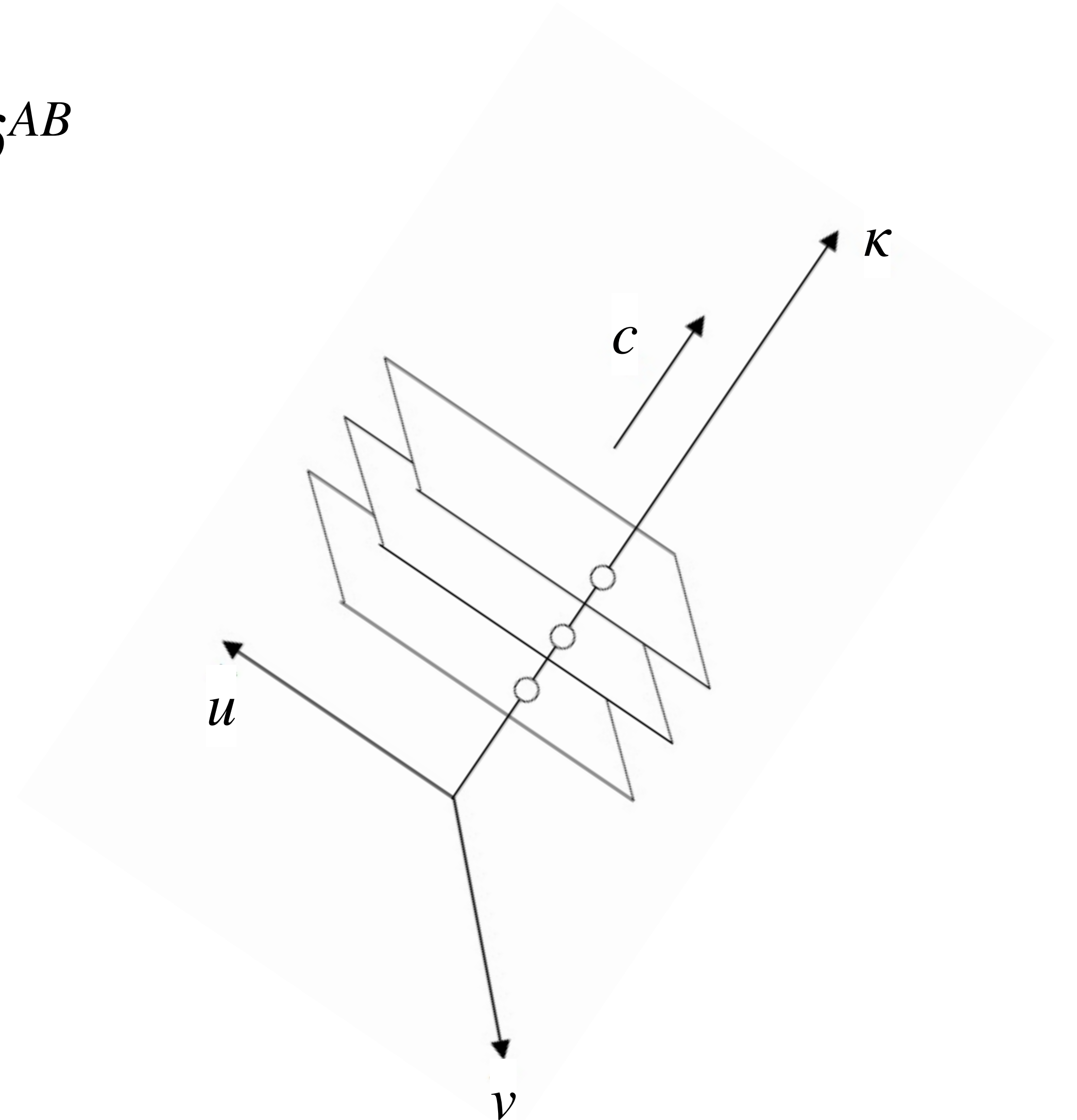
- $\tilde{h}_{ab} = \tilde{h}_+ e_{ab}^+ + \tilde{h}_\times e_{ab}^\times = \sum_{A=+,\times} \tilde{h}_A e_{ab}^A = \tilde{h}_A e_{ab}^A$



+ polarization

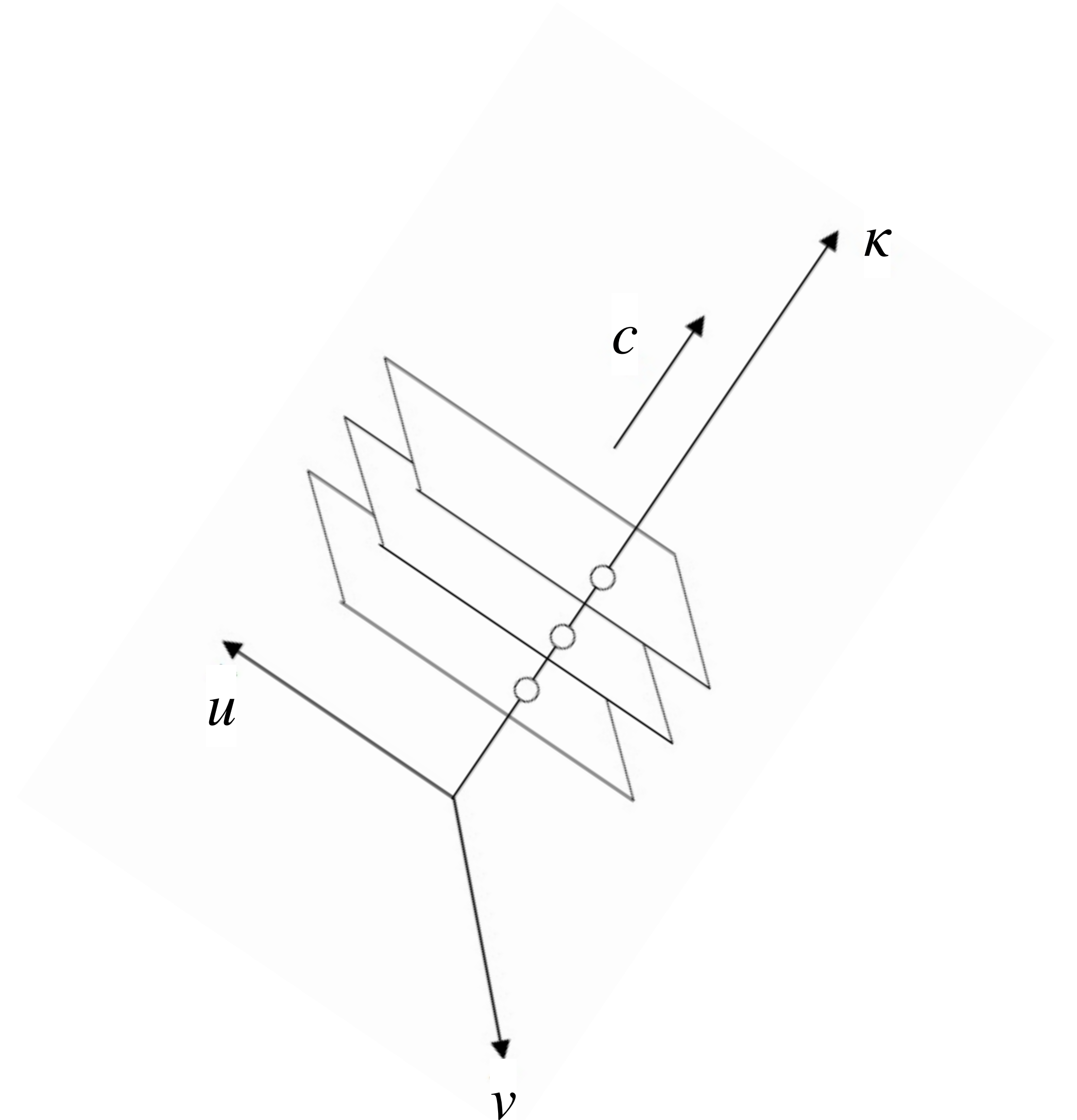


$\times$ polarization



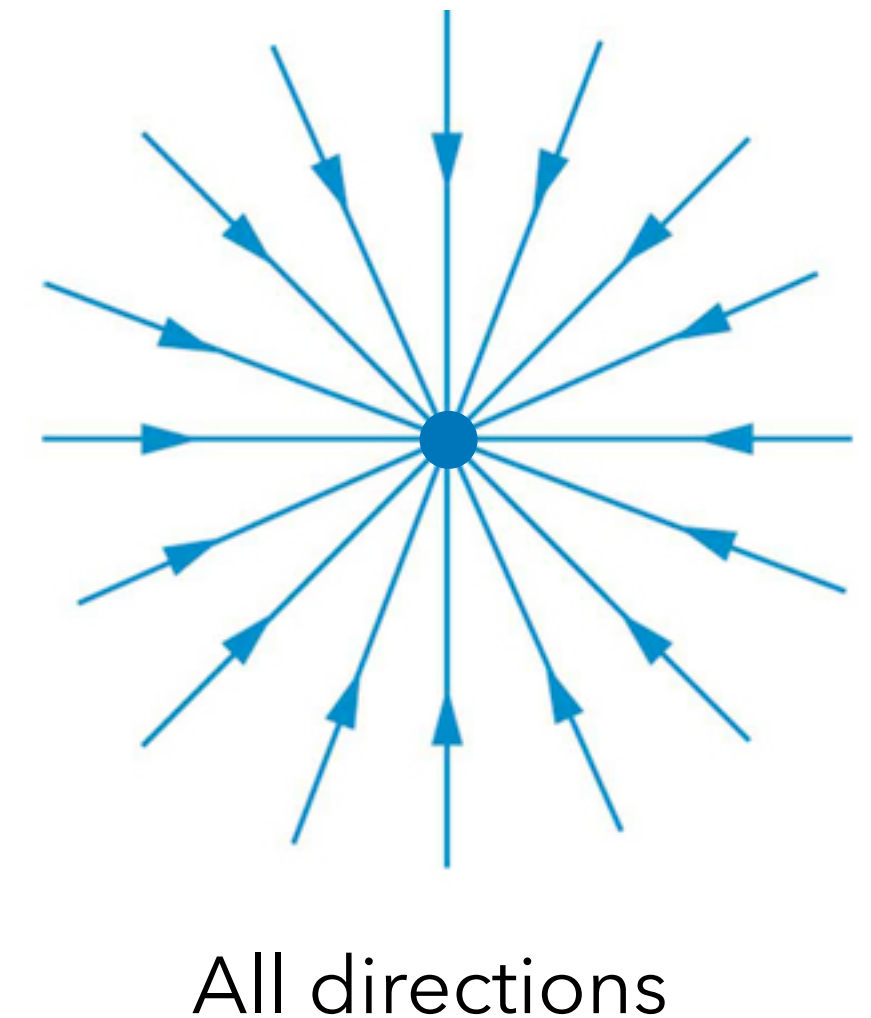
Plane GWs

- $h_{ab}(t, \vec{x}) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{h}_{ab}(\omega) e^{iP(t, \vec{x}; \omega)}$
- Complex amplitude: $\tilde{h}_{ab}(\omega) = \tilde{h}_A(\omega) e_{ab}^A$
- Note that $\tilde{h}_{ab}(-\omega) = \tilde{h}_{ab}^*(\omega)$
- Phase: $P(t, \vec{x}; \omega) = \omega(-t + \kappa \cdot \vec{x})$



General GWs

- $h_{ab}(t, \vec{x}) = \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{h}_{ab}(\omega, \kappa) e^{iP(t, \vec{x}; \omega, \kappa)}$
- Integration over all directions: $\int d^2\kappa = \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin \theta$
 - where $\kappa = \sin \theta \cos \phi \hat{x} + \sin \theta \sin \phi \hat{y} + \cos \theta \hat{z}$
- Complex amplitude: $\tilde{h}_{ab}(\omega, \kappa) = \tilde{h}_A(\omega, \kappa) e_{ab}^A(\kappa)$
- Note that $\tilde{h}_{ab}(-\omega, \kappa) = \tilde{h}_{ab}^*(\omega, \kappa)$
- Phase: $P(t, \vec{x}; \omega, \kappa) = \omega(-t + \kappa \cdot \vec{x})$
- All information of GWs are encoded in $\tilde{h}_{ab}(\omega, \kappa)$.



Stochastic GWs

- $h_{ab}(t, \vec{x}) = \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{h}_{ab}(\omega, \kappa) e^{iP(t, \vec{x}; \omega, \kappa)}$
- $\tilde{h}_{ab}(\omega, \kappa)$ is statistical random variable.
- First moment of random variable
 - $\langle \tilde{h}_{ab}(\omega, \kappa) \rangle = ?$
- Second moment of random variable
 - $\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega', \kappa') \rangle = ?$

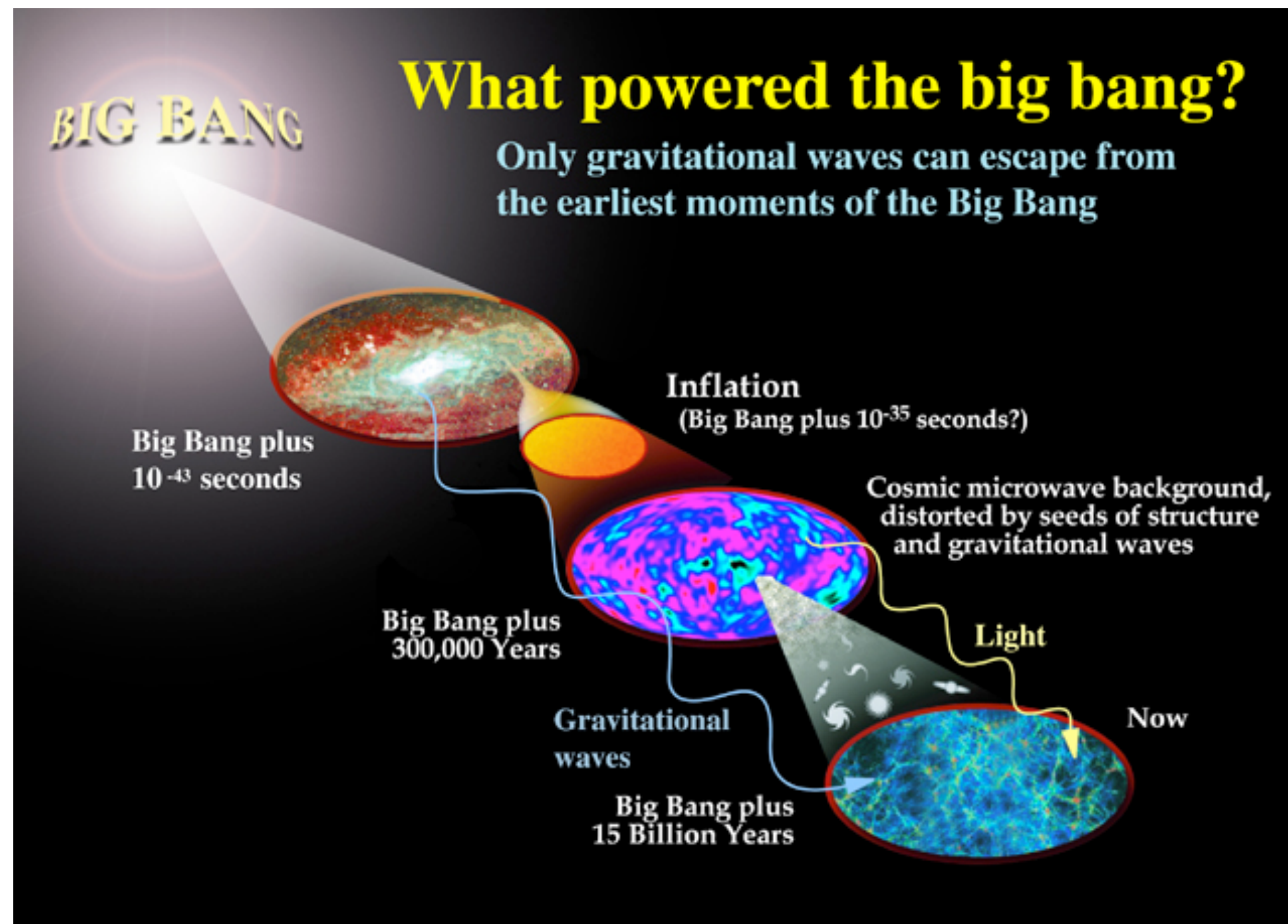


An example of stochastic process

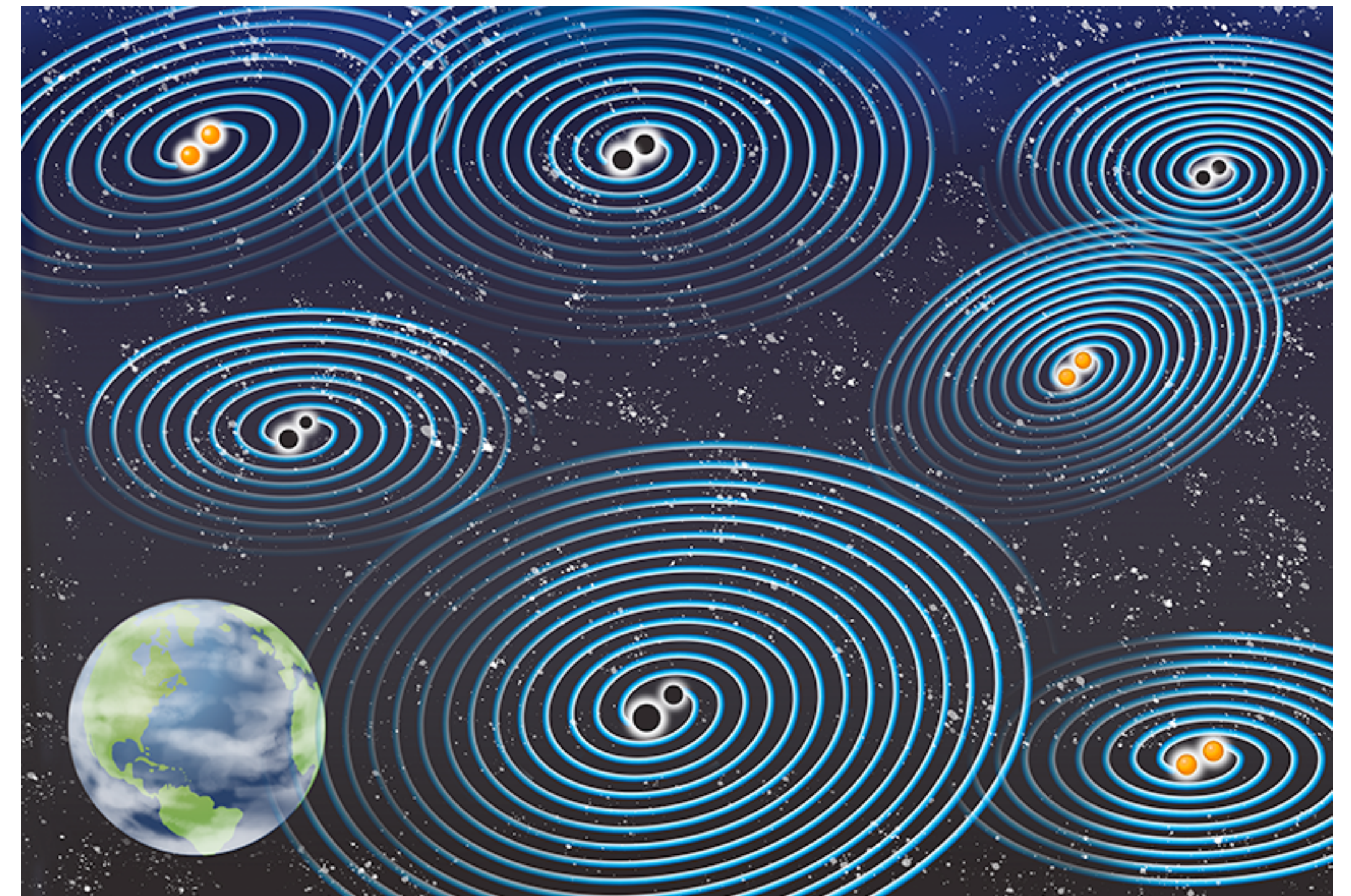


Stochastic Gravitational Wave Background

- Cosmological origin: Quantum state in early universe
- Astrophysical origin: Distribution of compact binaries



Stochastic GW from Big Bang / Credit: NASA



Distribution of Compact Binaries / Credit: APS

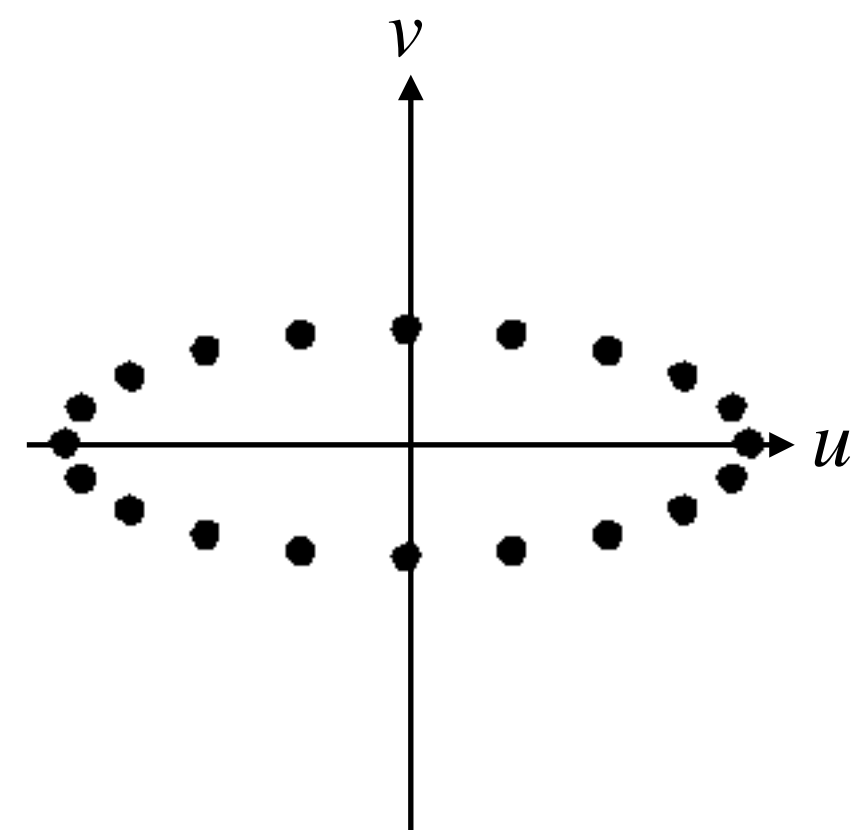
Statistical Assumptions on SGWB

- Gaussian and stationary assumptions
 - $\left\langle \tilde{h}_{ab}(\omega, \kappa) \right\rangle = 0$
 - $\left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega', \kappa') \right\rangle \propto S_h(\omega) \delta(\omega - \omega')$
 - $S_h(\omega)$: power spectral density (real and even)
- Isotropic assumption
 - $\left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega', \kappa') \right\rangle \propto \delta^2(\kappa - \kappa')$
 - $\delta^2(\kappa - \kappa') = \delta(\cos \theta - \cos \theta') \delta(\phi - \phi')$
 - $\kappa = \sin \theta \cos \phi \hat{x} + \sin \theta \sin \phi \hat{y} + \cos \theta \hat{z}$

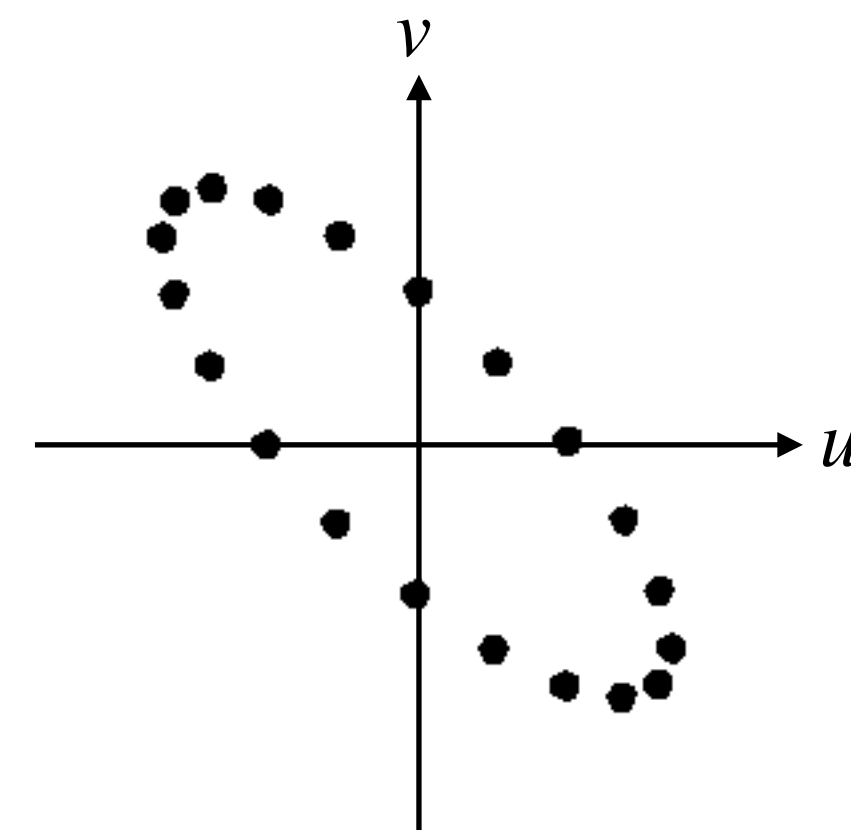


Statistical Assumptions on SGWB

- No prefer polarization assumption
 - $\left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega', \kappa') \right\rangle \propto \Lambda_{abcd}(\kappa)$
 - $\Lambda_{cd}^{ab} = e_A^{ab} e_{cd}^A = P_{(c}^a P_{d)}^b - \frac{1}{2} P^{ab} P_{cd}$: projection operator for symmetric traceless rank-2 tensors in $u - v$ plane
 - $P_b^a = \delta_b^a + n^a n_b - \kappa^a \kappa_b$: projection operator for vector to $u - v$ plane



+ polarization



x polarization

Second Moment of SGWBs

- $\left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega', \kappa') \right\rangle = 2\pi \delta(\omega - \omega') S_h(\omega) \frac{1}{4\pi} \delta^2(\kappa - \kappa') \frac{1}{2} \Lambda_{abcd}(\kappa)$

- All information of SGWB are encoded in $S_h(\omega)$

$$\begin{aligned} \langle h_{ab} h^{ab} \rangle &= \int d^2\kappa \int d^2\kappa' \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}^{ab*}(\omega', \kappa') \right\rangle e^{iP(t, \vec{x}; \omega, \kappa)} e^{-iP(t, \vec{x}; \omega', \kappa')} \\ &= \frac{1}{4\pi} \int d^2\kappa \int d^2\kappa' \delta^2(\kappa - \kappa') \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_h(\omega) e^{i\omega(-t + \kappa \cdot \vec{x})} e^{-i\omega(-t + \kappa' \cdot \vec{x})} \end{aligned}$$

- $$= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_h(\omega)$$

- It justifies the name of $S_h(\omega)$, i.e., the power spectral density.
- In usual text, the result has factor of 4 due to the careless normalization.

"Omega GW" of SGWB

- Energy density of SGWB

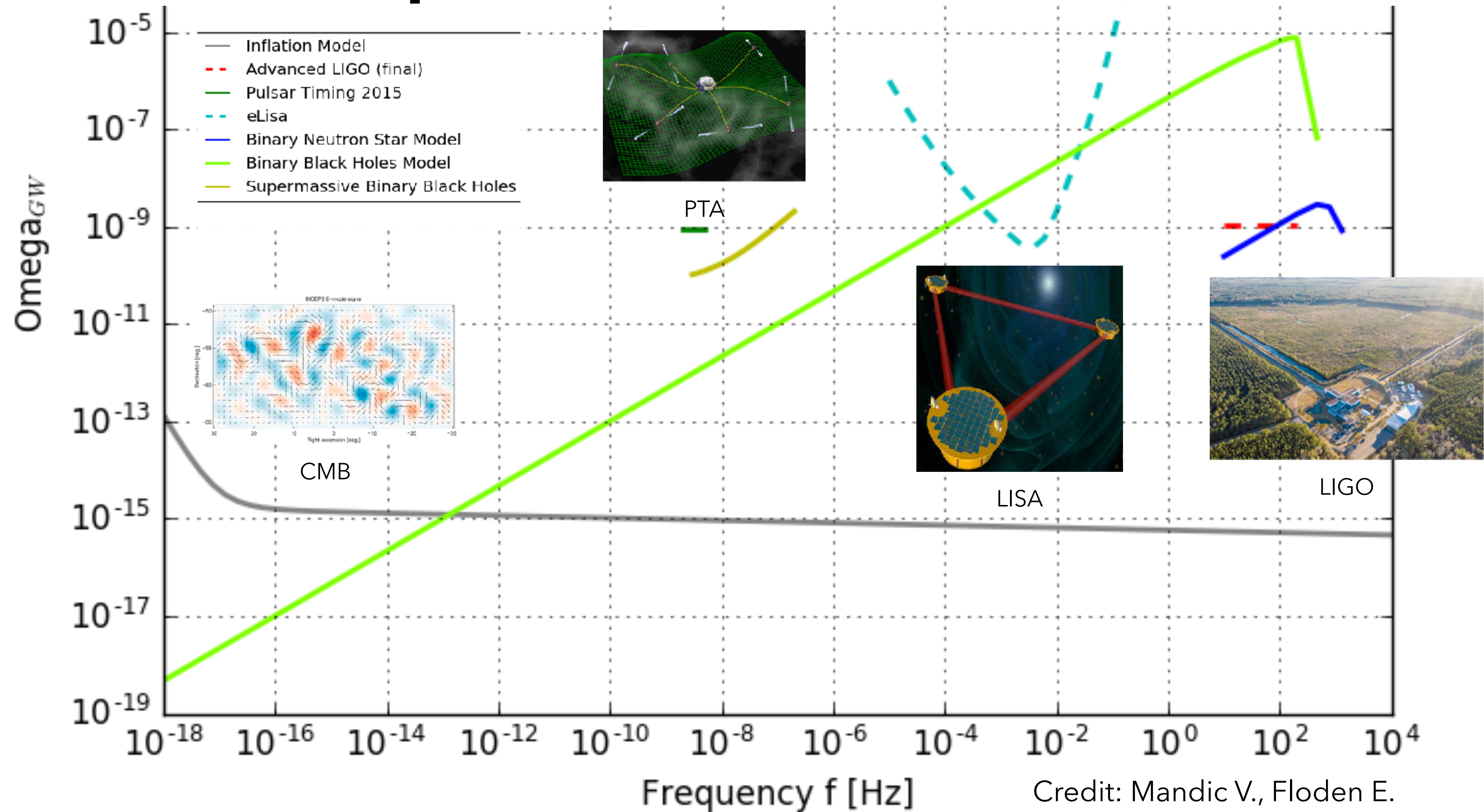
- $\rho_{gw} = \frac{1}{32\pi} \langle \partial_t h_{ab} \partial_t h^{ab} \rangle = \frac{1}{32\pi} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega^2 S_h(\omega) = \frac{1}{32\pi} \int_0^{\infty} df (2\pi f)^2 S_h^{\text{one-sided}}(f)$

- Omega GW

- $\Omega_{gw} = \frac{d(\rho_{gw}/\rho_c)}{d \ln f} = \frac{\pi^2}{3H_0^2} f^3 S_h^{\text{one-sided}}(f)$

- where $\rho_c = \frac{3H_0^2}{8\pi}$

Spectrum of SGWBs



Noise Reduction: Two-Detector Correlation Method

Output vs Signal

- $s(t) = h(t) + n(t)$

$$h(t) = \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{D}^{ab}(\omega) \tilde{h}_{ab}(\omega, \kappa) e^{i\omega(-t + \kappa \cdot \vec{x}_0)}$$

$$= \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{D}^{ab}(\omega) e_{ab}^A(\kappa) \tilde{h}_A(\omega, \kappa) e^{i\omega(-t + \kappa \cdot \vec{x}_0)}$$

- $$= \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{F}^A(\omega, \kappa) \tilde{h}_A(\omega, \kappa) e^{i\omega(-t + \kappa \cdot \vec{x}_0)}$$

- $\tilde{D}^{ab}$: detector tensor

- $\tilde{F}^A$: pattern function

- $$\tilde{h}(\omega) = \int d^2\kappa \tilde{D}^{ab}(\omega) \tilde{h}_{ab}(\omega, \kappa) e^{i\kappa \cdot \vec{x}_0}$$

Noise

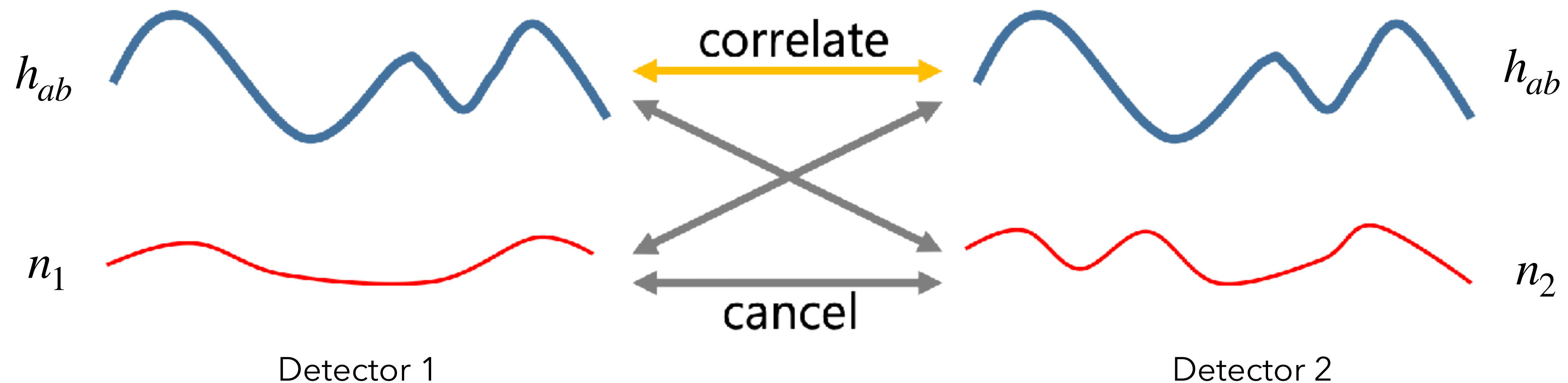
- $n(t) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{n}(\omega) e^{-i\omega t}$
- Gaussian and stationary assumptions
 - $\langle \tilde{n}(\omega) \rangle = 0$
 - $\langle \tilde{n}(\omega) \tilde{n}^*(\omega') \rangle = 2\pi \delta(\omega - \omega') S_n(\omega)$
 - $S_n(\omega)$: noise spectral density (real and even)
- All information of the noise are encoded in $S_n(\omega)$
- $\langle n^2(t) \rangle = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_n(\omega)$

Correlation Method

- Output for two-detectors

- $s_1(t) = h_1(t) + n_1(t) = \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{F}_1^A(\omega, \kappa) \tilde{h}_A(\omega, \kappa) e^{i(-t + \kappa \cdot \vec{x}_1)} + n_1(t)$

- $s_2(t) = h_2(t) + n_2(t) = \int d^2\kappa \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{F}_2^A(\omega, \kappa) \tilde{h}_A(\omega, \kappa) e^{i(-t + \kappa \cdot \vec{x}_2)} + n_2(t)$



Correlation Measure

$$\begin{aligned}
 Y &= \int_{-T/2}^{T/2} dt \int_{-T/2}^{T/2} dt' s_1(t) s_2(t') Q(t - t') \\
 &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega''}{2\pi} \tilde{s}_1(\omega) \tilde{s}_2^*(\omega') \tilde{Q}(\omega'') 2\pi\delta_T(\omega - \omega') 2\pi\delta_T(\omega' - \omega'')
 \end{aligned}$$

- $\simeq \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{s}_1(\omega) \tilde{s}_2^*(\omega) \tilde{Q}(\omega)$
- $Q(t)$: real filter function
- $\delta_T(\omega) = \frac{\sin(\omega T/2)}{\pi\omega} \rightarrow \delta(\omega)$ as $\omega T \rightarrow \infty$
- $\lim_{\omega \rightarrow 0} \delta_T(\omega) = T/2\pi$

Signal to Noise Ratio

- Signal: Expectation of Y

$$\text{Remark that } \left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega', \kappa') \right\rangle = 2\pi \delta(\omega - \omega') S_h(\omega) \frac{1}{4\pi} \delta^2(\kappa - \kappa') \frac{1}{2} \Lambda_{abcd}(\kappa)$$

- $S = \langle Y \rangle = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \left\langle \tilde{h}_1(\omega) \tilde{h}_2^*(\omega) \right\rangle \tilde{Q}(\omega)$

$$\left\langle \tilde{h}_1(\omega) \tilde{h}_2^*(\omega) \right\rangle = \int d^2\kappa \int d^2\kappa' \tilde{D}_1^{ab}(\omega) \tilde{D}_2^{*cd}(\omega) \left\langle \tilde{h}_{ab}(\omega, \kappa) \tilde{h}_{cd}^*(\omega, \kappa') \right\rangle e^{i\omega\kappa \cdot \vec{x}_1} e^{-i\omega\kappa' \cdot \vec{x}_2}$$

- $$\begin{aligned} &\simeq 2\pi \delta_T(0) S_h(\omega) \frac{1}{2} \tilde{D}_1^{ab}(\omega) \tilde{D}_2^{*cd}(\omega) \frac{1}{4\pi} \int d^2\kappa \Lambda_{abcd}(\kappa) e^{i\omega\kappa \cdot (\vec{x}_1 - \vec{x}_2)} \\ &= TS_h(\omega) \tilde{\Gamma}(\omega) \end{aligned}$$

- $\tilde{\Gamma}(\omega) = \frac{1}{2} \tilde{D}_1^{ab}(\omega) \tilde{D}_2^{*cd}(\omega) \frac{1}{4\pi} \int d^2\kappa \Lambda_{abcd}(\kappa) e^{i\omega\kappa \cdot (\vec{x}_1 - \vec{x}_2)}$: overlap reduction function

- $S = T \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_h(\omega) \tilde{\Gamma}(\omega) \tilde{Q}(\omega)$

Signal to Noise Ratio

- Noise: Standard deviation of Y without GWs

$$\begin{aligned}
 N^2 &= \left[\langle Y^2 \rangle - \langle Y \rangle^2 \right]_{h=0} \\
 &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \left\langle \tilde{n}_1(\omega) \tilde{n}_2^*(\omega) \tilde{n}_1^*(\omega') \tilde{n}_2(\omega') \right\rangle \tilde{Q}(\omega) \tilde{Q}^*(\omega') + \left[\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \left\langle \tilde{n}_1(\omega) \tilde{n}_2^*(\omega) \right\rangle \tilde{Q}(\omega) \right]^2 \\
 &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \left\langle \tilde{n}_1(\omega) \tilde{n}_1^*(\omega') \right\rangle \left\langle \tilde{n}_2(\omega') \tilde{n}_2^*(\omega) \right\rangle \tilde{Q}(\omega) \tilde{Q}^*(\omega') \\
 \bullet &= T \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_{n,1}(\omega) S_{n,2}(\omega) \left| \tilde{Q}(\omega) \right|^2 \\
 \bullet &\frac{S}{N} = \sqrt{T} \frac{\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_h(\omega) \tilde{\Gamma}(\omega) \tilde{Q}(\omega)}{\sqrt{\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_{n,1}(\omega) S_{n,2}(\omega) \left| \tilde{Q}(\omega) \right|^2}}
 \end{aligned}$$

Maximal Signal to Noise Ratio

- We have to determine $Q(t)$ to maximize SNR.

- $$\frac{S}{N} = \sqrt{T} \frac{\langle \tilde{Q}, \tilde{\Gamma} S_h / S_{n,1} S_{n,2} \rangle}{\sqrt{\langle \tilde{Q}, \tilde{Q} \rangle}}$$

- where $\langle \tilde{A}, \tilde{B} \rangle = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{A}(\omega) \tilde{B}^*(\omega) S_{n,1}(\omega) S_{n,2}(\omega)$

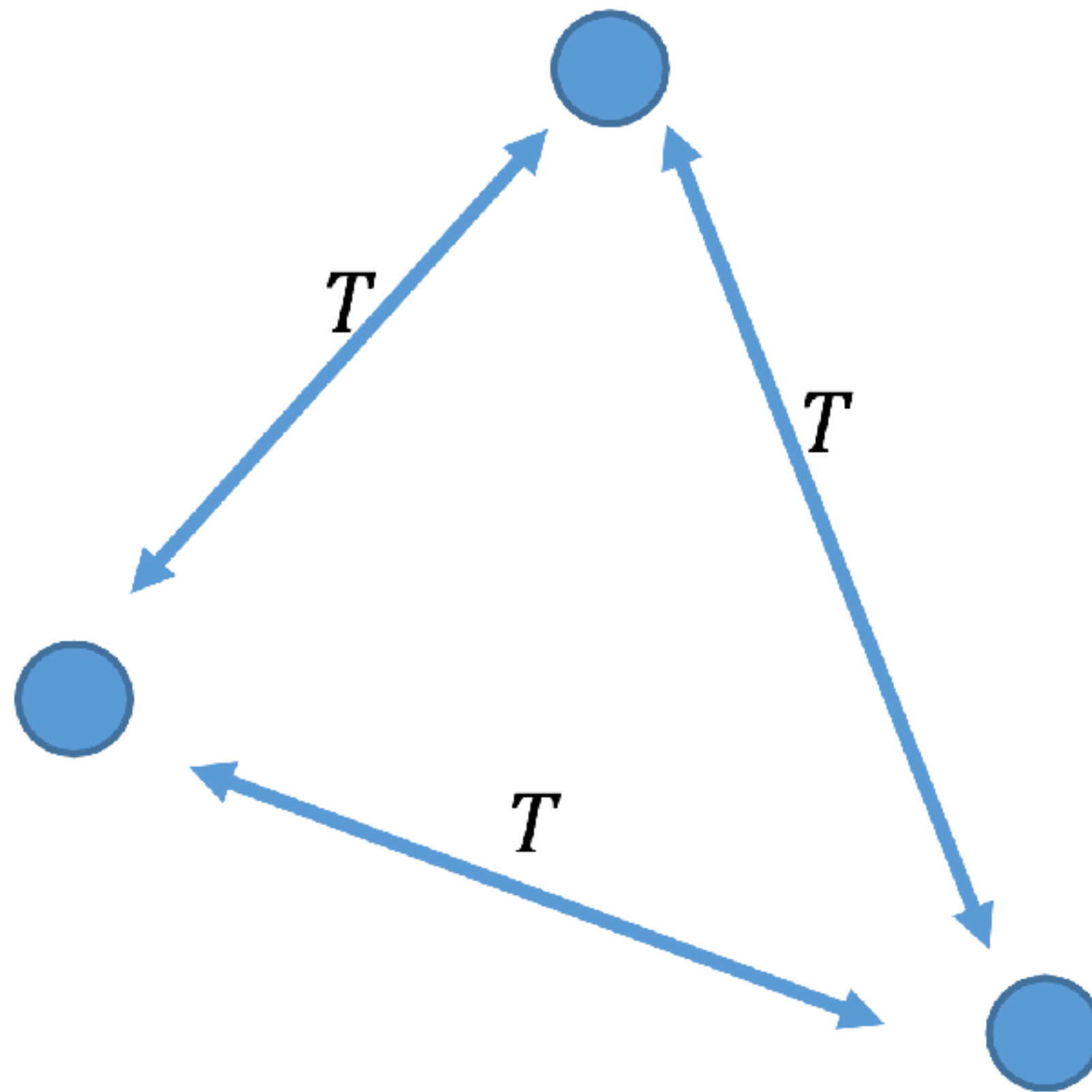
- SNR is maximized when $\tilde{Q} = \tilde{\Gamma} S_h / S_{1,n} S_{2,n}$

- Then, maximal SNR is

- $$\frac{S}{N} = \sqrt{T} \left[\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \left| \tilde{\Gamma}(\omega) \right|^2 \frac{S_h^2(\omega)}{S_{n,1}(\omega) S_{n,2}(\omega)} \right]^{1/2}$$

N-detector network

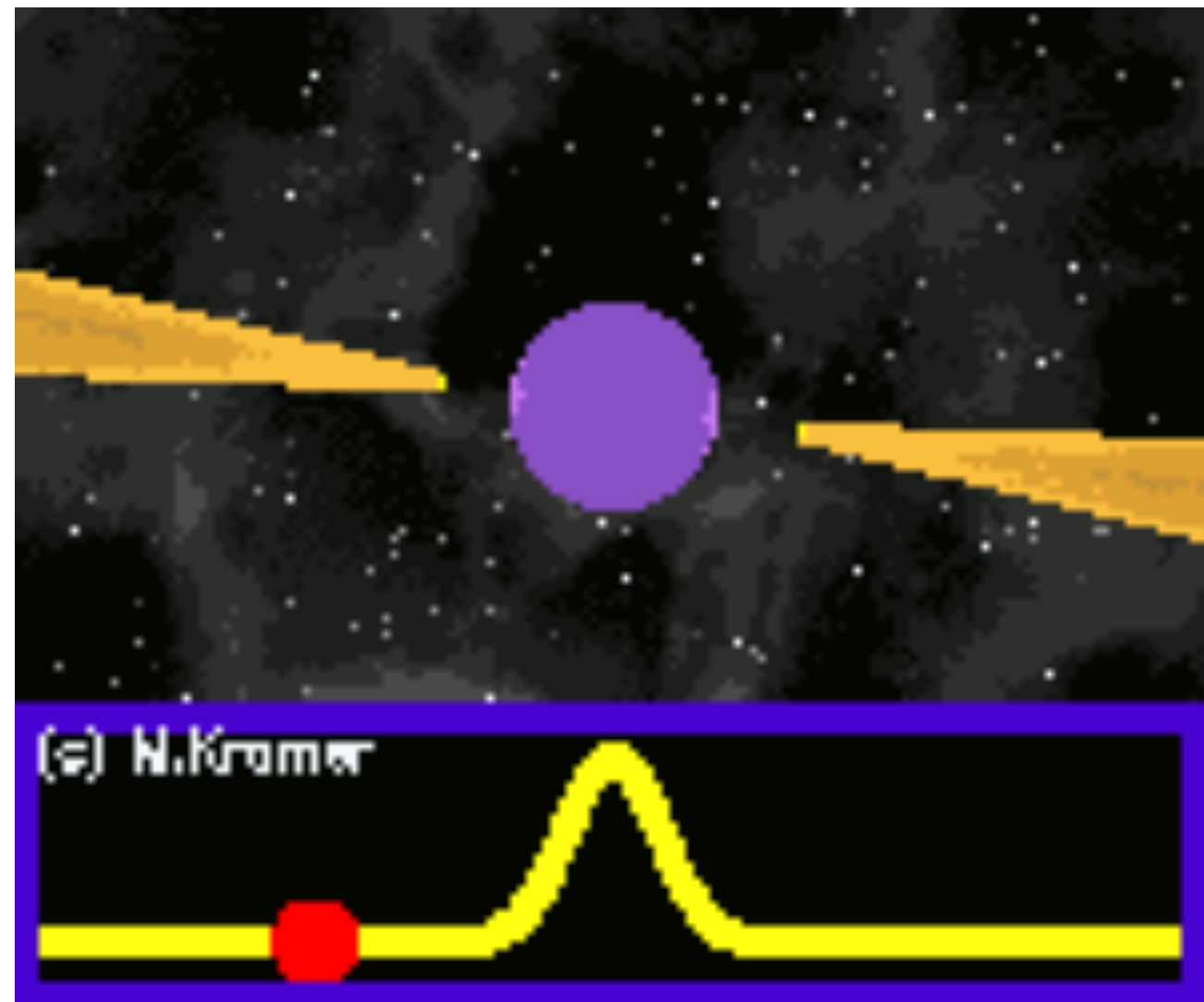
- Total measurement time: $T_N = \frac{N(N-1)}{2}T$



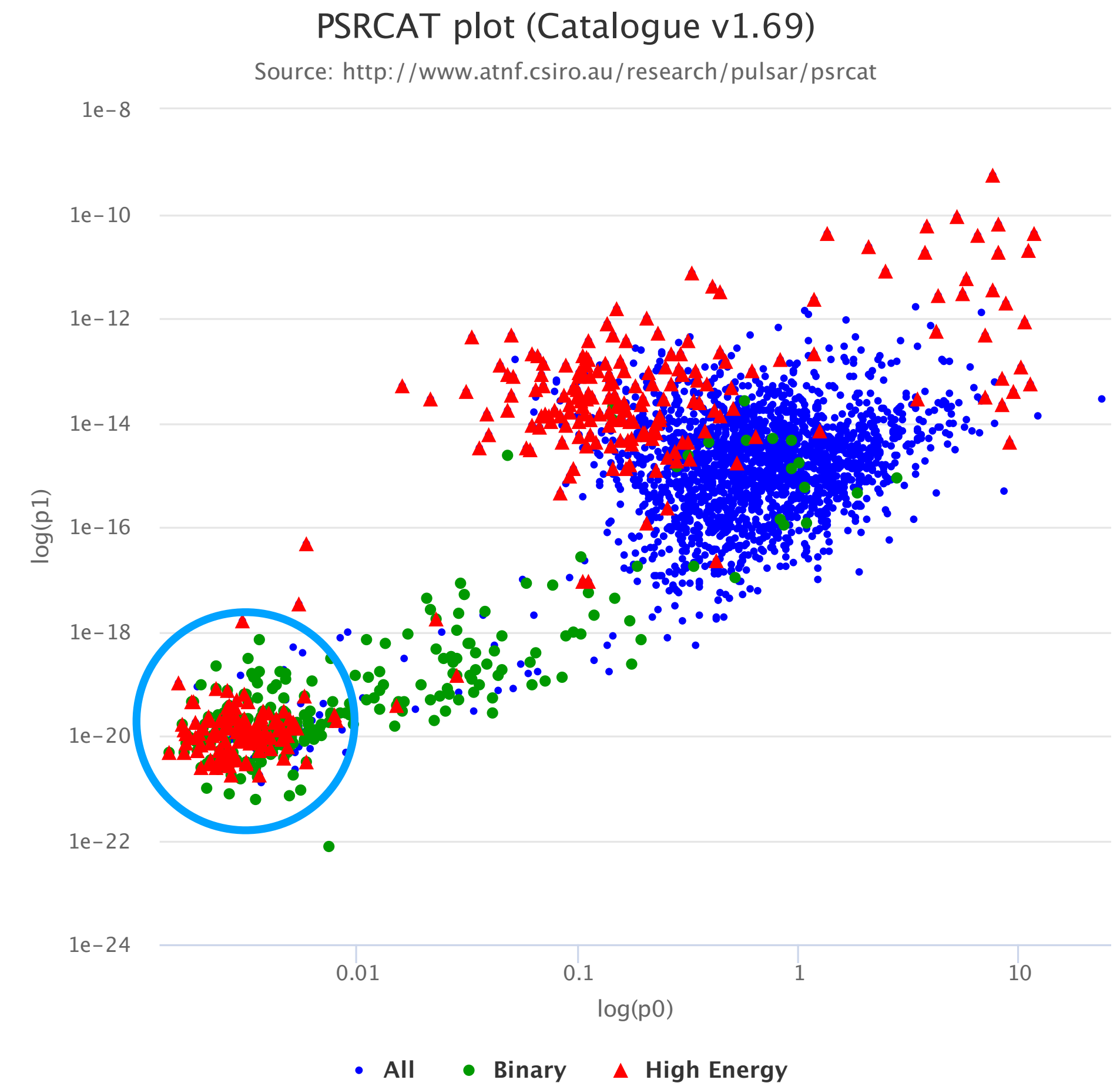
Pulsar Timing Arrays

Millisecond Pulsars (MSPs)

- Rotational period < 10 ms
- Extremely stable rotation
- Usually recycled by a companion star



Credit: Michael Kramer (JBCA,
University of Manchester).



Highcharts.com

ATNF Pulsar Catalogue

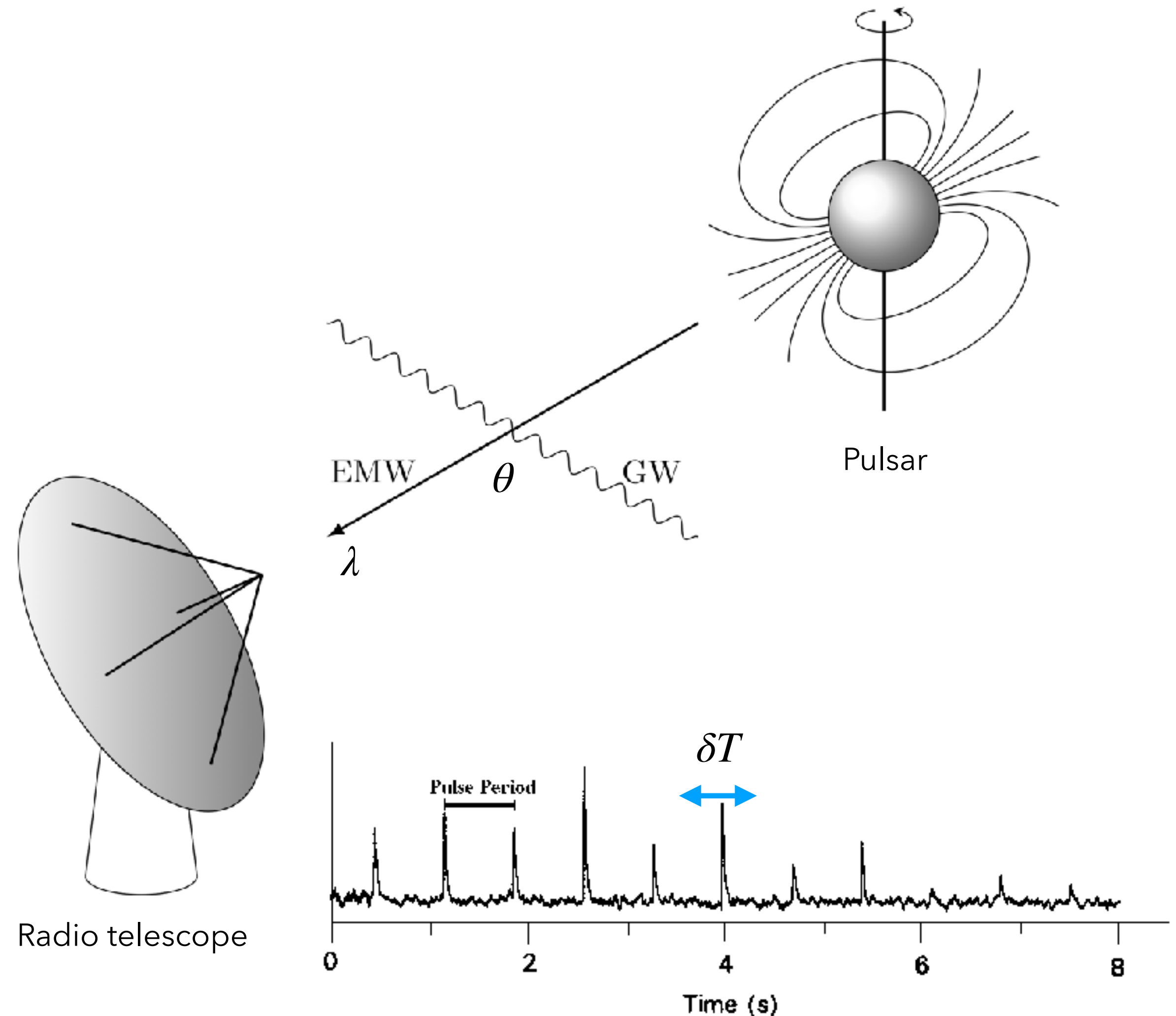
Perturbation of the Period by GWs

- Electromagnetic waves (EMWs) are perturbed in a spacetime with gravitational waves (GWs).

- Fractional change of the period

$$\frac{\delta T}{T} = \int d^2\kappa \int \frac{d\omega}{2\pi} \left[-\frac{\tilde{h}_{ab}\lambda^a\lambda^b}{2(1 - \cos\theta)} (1 - e^{i\Delta}) e^{iP} \right]$$

- $\tilde{h}$: complex amplitude of GW
- λ : spatial unit vector of EMW propagation
- θ : angle between EMW and GW
- Δ : time delay phase at pulsar
- P : phase of GW



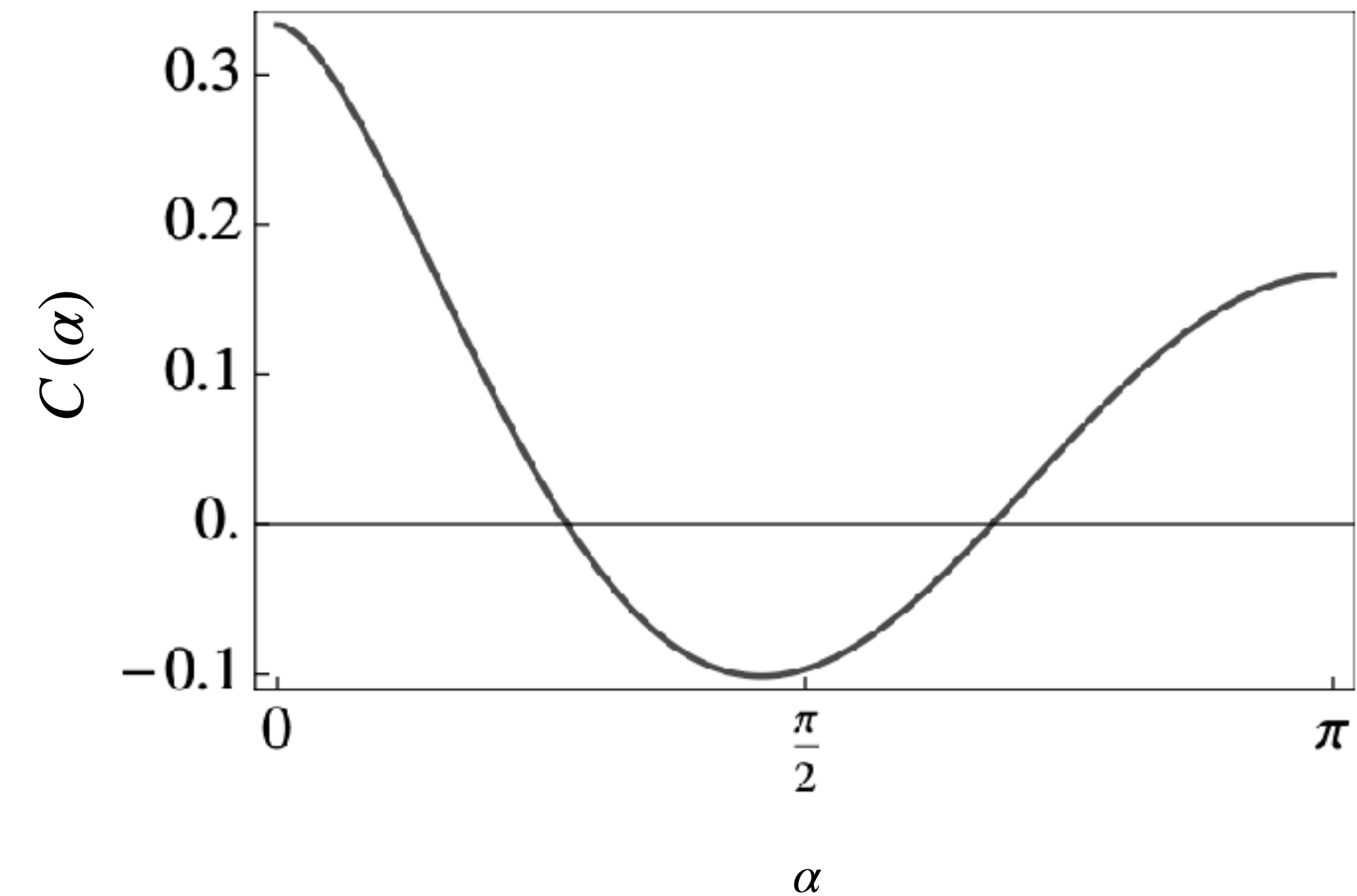
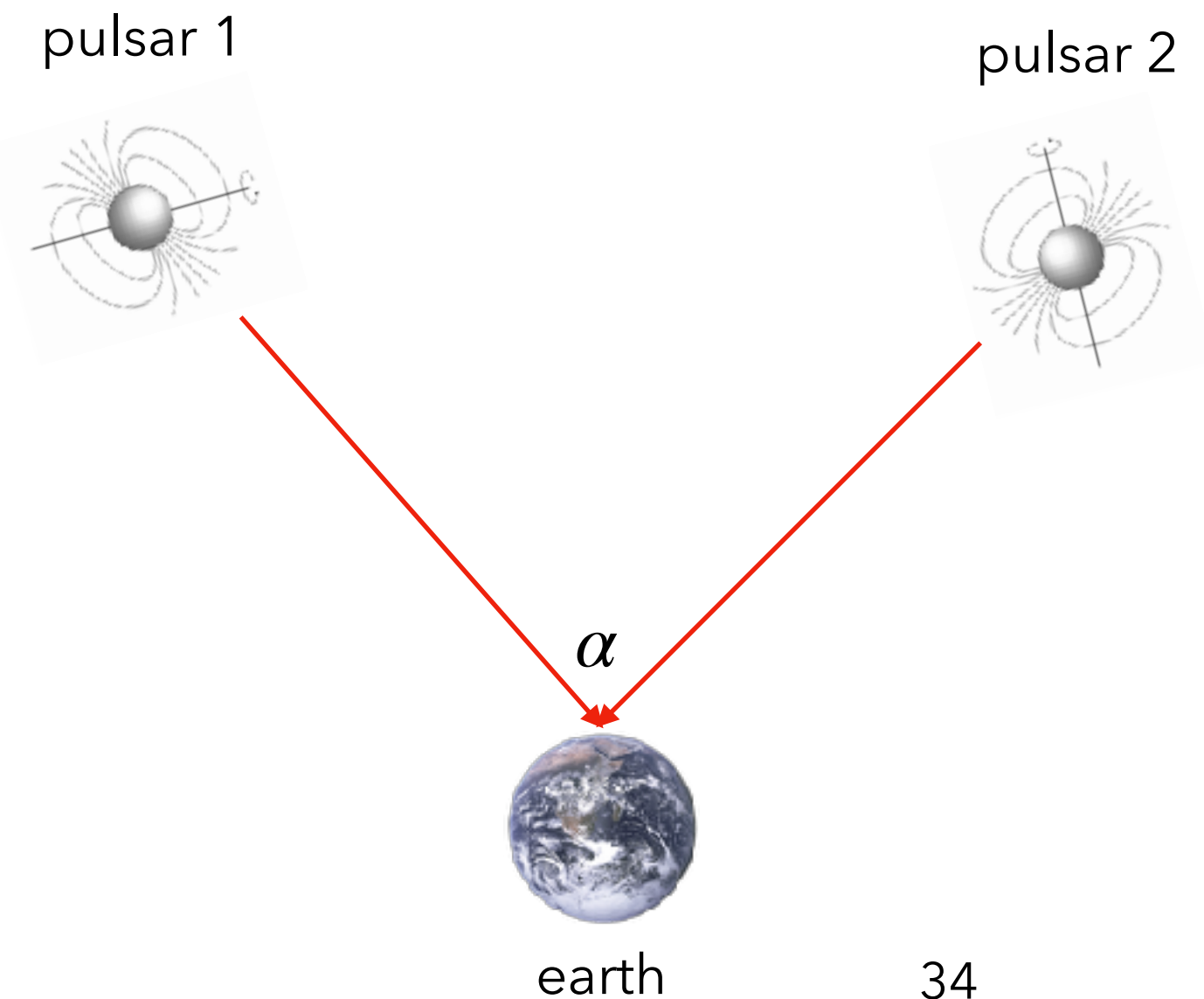
Pulsar Timing Array (PTA)

- Power of SGWB

- $\langle h_{ab} h^{ab} \rangle = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_h(\omega)$

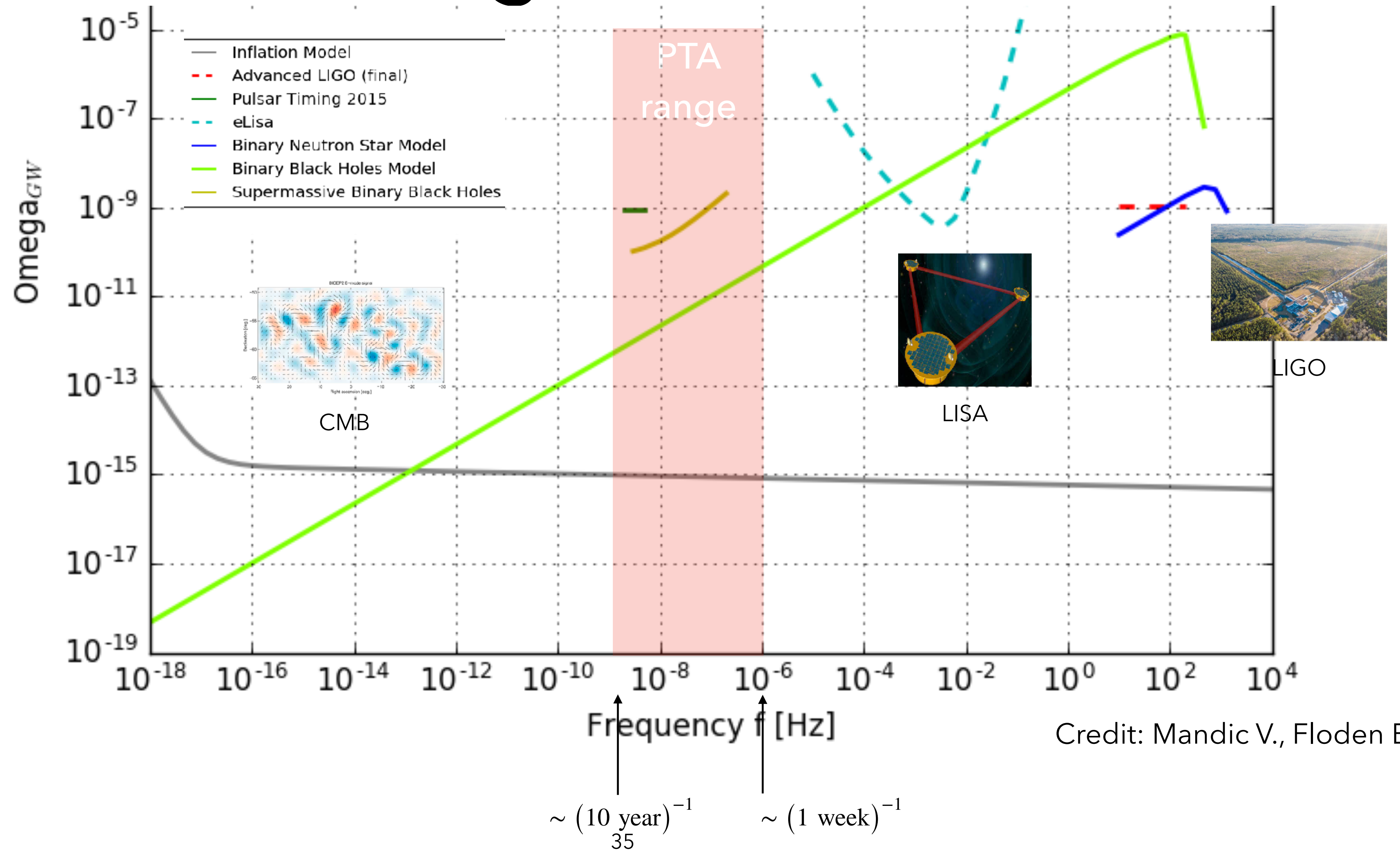
- Correlation of measurements from two pulsars

- $\left\langle \frac{\delta T_1}{T_1} \frac{\delta T_2}{T_2} \right\rangle = C(\alpha) \langle h_{ab} h^{ab} \rangle$



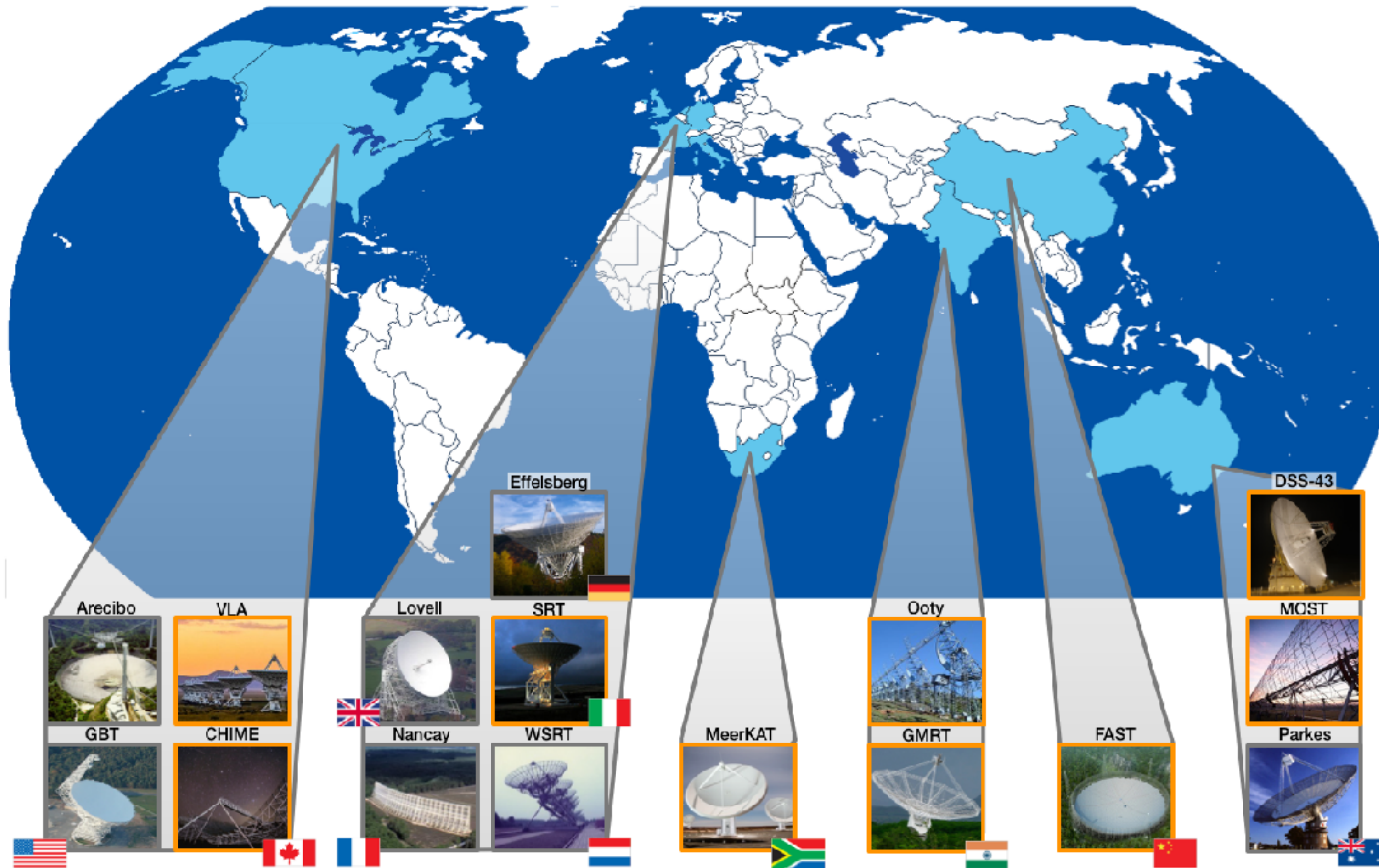
Hellings and Downs / ApJ (1983)

Target SGWB



IPTA

- EPTA+InPTA+NanoGrav+PPTA



Credit: NanoGrav



Square Kilometer Array Observatory (SKAO)

- SKA Low: 50 - 350 MHz / 131,072 antennas
- SKA Mid: 350 MHz - 15.4 GHz / 197 dishes
- Start from 2028



SKA Low / Australia



SKA Mid / South Africa

Summary

- SGWB has two origin. (cosmological / astrophysical)
- PTA targets SGWB comes from supermassive binaries with the period ~ 10 years.
- PTAs has accumulated data for ~ 15 years.
- We will detect SGWB soon.